

Rotational spectroscopy as a tool to investigate interactions between vibrational polyads in symmetric top molecules: low-lying states $\nu_8 \leq 2$ of methyl cyanide, CH_3CN

Holger S.P. Müller^{a,*}, Linda R. Brown^b, Brian J. Drouin^b, John C. Pearson^b, Isabelle Kleiner^c, Robert L. Sams^d, Keeyoon Sung^b, Matthias H. Ordu^a, Frank Lewen^a

^a*I. Physikalisches Institut, Universität zu Köln, Zùlpicher Str. 77, 50937 Köln, Germany*

^b*Jet Propulsion Laboratory, California Institute of Technology, Pasadena, CA 91109-8099, USA*

^c*Laboratoire Interuniversitaire des Systèmes Atmosphériques (LISA), UMR CNRS 7583, Universités Paris Est Créteil et Paris Diderot, Institut Pierre Simon Laplace, 61 Avenue du Général de Gaulle, 94010 Créteil Cedex, France*

^d*Pacific Northwest National Laboratory, P.O. Box 999, Mail Stop K8-88, Richland, WA 99352, USA*

Abstract

Rotational and rovibrational spectra of methyl cyanide were recorded to analyze interactions in low-lying vibrational states and to construct line lists for radio astronomical observations in space as well as for infrared spectroscopic investigations of planetary atmospheres. The rotational spectra cover large portions of the 36–1627 GHz region. In the infrared (IR), a spectrum was recorded for this study in the region of $2\nu_8$ around 717 cm^{-1} with assignments covering $684\text{--}765\text{ cm}^{-1}$. Additional spectra in the ν_8 region were used to validate the analysis.

Information on the K level structure of CH_3CN is almost exclusively obtained from IR spectra, as are basics of the J level structure. The large amount and the high accuracy of the rotational data improves knowledge of the J level structure considerably. Moreover, since these data extend to much higher J and K quantum numbers, they allowed us to investigate for the first time in depth local interactions between these states which occur at high K values. In particular, we have detected several interactions between $\nu_8 = 1$ and 2. Notably, there is a strong $\Delta\nu_8 = \pm 1$, $\Delta K = 0$, $\Delta l = \pm 3$ Fermi resonance between $\nu_8 = 1^{-1}$ and $\nu_8 = 2^{+2}$ at $K = 14$. Pronounced effects in the spectrum are also caused by resonant $\Delta\nu_8 = \pm 1$, $\Delta K = \mp 2$, $\Delta l = \pm 1$ interactions between $\nu_8 = 1$ and 2 at $K = 13$, $l = -1/K = 11$, $l = 0$ and at $K = 15$, $l = +1/K = 13$, $l = +2$. An equivalent resonant interaction occurs between $K = 14$ of the ground vibrational state and $K = 12$, $l = +1$ of $\nu_8 = 1$ for which we present the first detailed account. A preliminary account was given in an earlier study on the ground vibrational state. Similar resonances were found for CH_3CCH and, more recently, for CH_3NC , warranting comparison of the results. From data pertaining to $\nu_8 = 2$, we also investigated rotational interactions with $\nu_4 = 1$ as well as $\Delta\nu_8 = \pm 1$, $\Delta K = 0$, $\Delta l = \pm 3$ Fermi interactions between $\nu_8 = 2$ and 3.

We have derived N_2 - and self-broadening coefficients for the ν_8 , $2\nu_8 - \nu_8$, and $2\nu_8$ bands from previously determined ν_4 values. Subsequently, we determined transition moments and intensities for the three IR bands.

Keywords:

rotational spectroscopy, infrared spectroscopy, vibration-rotation interaction, methyl cyanide, interstellar molecule

1. Introduction

Among the more than 180 different molecules detected in the interstellar medium (ISM) or in circumstellar envelopes (CSEs) of late type stars, there are 56 with a CN moiety, see, e.g., the Molecules in Space page¹ of the Cologne Database for Molecular Spectroscopy, CDMS [1, 2] or the Bibliography of Astromolecules² of the Astrochymist web-page³. Almost two thirds of these can be viewed as cyanides in a wider sense. The alkyl cyanides n -propyl cyanide, $n\text{-C}_3\text{H}_7\text{CN}$, also known as 1-cyanopropane or butyronitrile, and iso -propyl cyanide, $i\text{-C}_3\text{H}_7\text{CN}$, also known as 2-cyanopropane or iso -butyronitrile,

both recently detected in the very massive and luminous galactic center source Sagittarius B2(N) (Sgr B2(N) for short) [3, 4] are the largest complex organic molecules detected in the ISM thus far, and, after HC_{11}N , the second largest cyanides. Moreover, iso -propyl cyanide was the first branched alkyl compound detected in space [4]. Ethyl cyanide, $\text{C}_2\text{H}_5\text{CN}$, the next smaller homolog, was detected much earlier in the Orion Molecular Cloud as well as Sgr B2 [5]. Its ^{13}C isotopologs were detected in Orion IRC2 [6] and Sgr B2(N) [7]. In addition, rotational transitions pertaining to the two lowest-lying vibrational states were detected in G327.3–0.6 [8] and Sgr B2(N) [9]. Transitions of higher excited states were detected in molecular clouds of Orion [10] and Sgr B2(N) [11].

Methyl cyanide, CH_3CN , also known as acetonitrile or cyanomethane, is the smallest alkyl cyanide. It is a trace constituent in Earth's atmosphere which was studied, e.g., by high-resolution spectroscopy using microwave limb sounding [12]

*Corresponding author.

Email address: hspm@ph1.uni-koeln.de (Holger S.P. Müller)

¹<https://cdms.astro.uni-koeln.de/classic/molecules>

²http://www.astrochymist.org/astrochymist_mole.html

³<http://www.astrochymist.org/>

and balloon-borne IR spectroscopy [13]. It is produced mainly by biomass burning [14, 15]. Acetonitrile is also a useful solvent in the chemical industry.

Methyl cyanide is a very important molecule in space. First identified in Sgr A and B more than 40 years ago, it was among the first molecules to be observed by radio astronomy [16]. It was also seen in comets, such as Kohoutek [17], in the atmosphere of Saturn’s moon Titan [18], in dark clouds, such as TMC-1 [19], around the low-mass protostar IRAS 16293–2422 [20], in the CSE of the famous carbon-rich late-type star CW Leonis, also known as IRC +10216 [21], and in external galaxies, such as NGC 253 [22]. Several isotopologs of CH₃CN were detected as well in the ISM, including the rare CH₂DCN [23] and possibly even ¹³CH₃¹³CN [11]. Transitions of methyl cyanide within its lowest excited vibrational state $v_8 = 1$ were detected more than 30 years ago [24]. More recently, brief accounts on observations of transitions belonging to $v_8 = 2$ [11, 25], $v_4 = 1$ [11] as well as on transitions of vibrationally excited methyl cyanide with ¹³C [11] have been published.

The identifications were based on extensive laboratory investigations which began almost 70 years ago [26]. The most recent account on the ground state rotational spectrum of the main isotopolog CH₃CN was given by some of the present authors along with analyses for several minor isotopologs [27]; unlabeled atoms refer to ¹H, ¹²C, and ¹⁴N. The CH₃CN line list was based mainly on an accurate and extensive study [28] which employed earlier data for frequencies below 74 GHz [29, 30, 31]. Information on the purely K -dependent terms cannot be determined by rotational spectroscopy for a strongly prolate symmetric top, such as methyl cyanide, unless from the very weak $\Delta K = 3$ transitions that gain intensity from centrifugal distortion effects on the the dipole moment. Thus the purely K -dependent terms A , D_K , and an estimate of H_K were determined from $\Delta K = 3$ ground state energy differences [32].

The lowest excited vibrational state of CH₃CN is the doubly degenerate CCN bending mode $v_8 = 1$ at a vibrational energy of 365.0 cm⁻¹ [33]. It was studied first in 1961 by rotational spectroscopy between 92 and 222 GHz ($J'' = 4$ to 11) [34]. Soon thereafter, more accurate data with partly resolved ¹⁴N hyperfine structure between 18 and 148 GHz ($J'' = 0$ to 7) were reported [35, 36]. Later, data were obtained near 370 GHz ($J'' = 19$) [37] and, employing a laser-sideband spectrometer, in selected regions between 589 and 1395 GHz ($J'' = 31$ to 75) [38]. Direct l -type transitions between the $k \times l = +1$ components were also observed with $21 \leq J \leq 45$ [39, 40].

To our knowledge, there was only one prior high-resolution IR study of the v_8 band [33], although there were earlier low-to moderate-resolution studies of v_8 or indirect determinations of the vibrational energy from higher vibrational bands. A resonant interaction with $\Delta k = -2$ and $\Delta l = +1$ between $v = 0$, $K = 14$ and $v_8 = 1^{+1}$, $K = 12$ at $J = 43$ was reported [41], but could only be analyzed approximately because only two slightly perturbed transitions in the ground vibrational state ($J'' = 38$ and 39) were observed.

The parallel $2v_8$ band at 716.7 cm⁻¹ [42] was treated as an isolated state in that study even though a rather strong Fermi

interaction of $v_8 = 2^0$ with $v_4 = 1$, the non-degenerate C–C stretching mode at 920.3 cm⁻¹ ($W_{488} \approx 30$ cm⁻¹) was proposed from a force field analysis [43]. This interaction has not been considered in high resolution analyses of $v_8 = 2$ thus far. Transitions to $v_8 = 2^{\pm 2}$ became observable only through q_{22} interaction with $v_8 = 2^0$ [42]. A Fermi interaction between $v_7 = 1^{+1}$, and $v_8 = 2^{-2}$ at $K = 12$ and 13 was proposed to explain deviations in the high- K part of the v_7 band [44]. This interaction was confirmed in a later study mainly concerned with v_4 , v_7 and $3v_8^1$ bands between 850 and 1150 cm⁻¹ [45]; in addition, a $\Delta k = -1$ and $\Delta l = +2$ interaction between $v_8 = 2^{-2}$ and $v_4 = 1$ at $K = 6$ and 5, respectively was found for the first time along with several other perturbations among the higher lying states.

The rotational spectrum of CH₃CN in its $v_8 = 2$ state has also been studied between 18 and 112 GHz ($J'' = 0$ to 5) [36, 46, 47], near 370 GHz ($J'' = 19$) [37] and, employing a laser-sideband spectrometer, in selected regions between 787 and 1273 GHz ($J'' = 20$ to 68) [38].

Despite the numerous studies of methyl cyanide in various spectral regions, there has been no detailed attempt to study its spectroscopy globally. If there were effects, which were caused by vibrational states outside of the polyad under investigation, they were only considered through the effective parameters, see, e.g., Refs. [41, 45].

In the present work, we have recorded extensive rotational and rovibrational spectra in order to study vibration-rotation interactions not only within vibrational states or within vibrational polyads, but also between vibrational polyads. In comparison to earlier data, we access higher J and K quantum numbers, have a more complete coverage of the data and typically higher accuracies. The data permitted the first detailed analysis of a resonant interaction between $v = 0$ and $v_8 = 1$ and the first detections and analyses of interactions between $v_8 = 1$ and 2 as well as between $v_8 = 2$ and 3. Preliminary results were described earlier [48]. In this article, we present an expanded analysis of these vibrational states of methyl cyanide with v_8 up to 2 and provide preliminary analyses of interactions with vibrational states pertaining to the next polyad. We calculate rotational and rovibrational spectra from our resulting analysis to support methyl cyanide observations in diverse environments in space as well as in Earth’s atmosphere and those of other planets or related objects.

Interactions between vibrational polyads of symmetric top molecules, such as methyl cyanide, are facilitated by the comparatively small vibrational energy of 365 cm⁻¹ for $v_8 = 1$ combined with the large Coriolis coupling between the two l components which causes levels with $\Delta K = 0$, but different l to repel each other increasingly at higher K such that the lower energy l component can interact with the ground vibrational state at specific K whereas the higher energy l component in turn can interact with the different l components of $v_8 = 2$ also at specific K . Vibrational polyads are groups of one or more vibrational states separated by approximately equal amounts of vibrational energy, see e.g. Ref. [49] for a recent review.

Interactions in symmetric top molecules between vibrational polyads were studied first in the case of propyne, CH₃CCH, also known as methylacetylene, which is isoelectronic to methyl

cyanide. A $\Delta v = \pm 1$, $\Delta K = 0$, $\Delta l = \pm 3$ Fermi interaction at $K = 12$ between $v_{10} = 1^{-1}$ and $v_{10} = 2^{+2}$ as well as a Coriolis interaction with $\Delta K \mp 1$, $\Delta l = \pm 2$ at $K = 10$ and 9 , respectively, between $v_{10} = 1^{-1}$ and $v_9 = 1^{+1}$ were first revealed mainly by rotational spectroscopy [50] and later improved by infrared as well as additional rotational spectroscopy [51]. The latter study also revealed a rotational (Coriolis-type) resonance between $v_{10} = 2$ and $v_5 = 1$, which is the lowest vibration of the next higher tetrad. Interactions between the dyad and the tetrad were studied in greater detail later [52]. More recently, such Fermi and α -resonances were analyzed between $v_8 = 1$ and 2 in methyl isocyanide [53]. The vibrational energy of $v_8 = 1$ is lower, only 267.3 cm^{-1} , such that the Fermi interaction occurs already at $K = 9$ and 10 , while they occur at $K = 12$ and 14 for the corresponding resonances in methylacetylene and methyl cyanide.

We are not aware of resonant interactions between the ground and an excited vibrational state in the case of a symmetric top molecule other than the case of CH_3CN . However, a Coriolis-type interaction between the ground vibrational state of HNCO and its lowest $v_5 = 1$ state was proposed and its strength estimated [54]. An analysis was presented for the analogous case of the isoelectronic HN_3 molecule [55]. A remarkable case is vinyl cyanide, $\text{C}_2\text{H}_3\text{CN}$, also known as acrylonitrile, for which the ground vibrational state was connected to its lowest $v_{11} = 1$ state through $\Delta K = 5$ and 6 interactions without observation of high-resolution infrared transitions [56]. Recently, $v_{11} = 1$ was connected to $v_{15} = 1$ and $v_{11} = 2$ as well as the latter two to each other [57]. The density of vibrational states increases with energy, and therefore chances to connect states at lower quantum numbers and through interactions with smaller ΔK increase. However, observation may be difficult in practice because transitions in higher vibrational states are weaker than corresponding ones in lower vibrational states. In addition, the number of at least similarly strong rotational transitions per frequency unit increases also with vibrational state.

2. Experimental details

2.1. Rotational spectra at the Jet Propulsion Laboratory

Most of the CH_3CN rotational data were extracted from broad frequency scans taken with the JPL cascaded multiplier spectrometer [58]. Generally, a multiplier chain source is passed through a $1 - 2$ meter pathlength flow cell and is detected by a silicon bolometer cooled to near 1.7 K . The cell is filled with a steady flow of reagent grade acetonitrile at room temperature, and the pressure and modulation are optimized to enable good S/N ratios with narrow lineshapes. The S/N ratio was optimized for a higher- K transition (e.g. $K = 12$) because of the very strong ground state transitions of the main isotopolog with lower K , which may exhibit saturated line profiles. This procedure enables better dynamic range for the extraction of line positions for rare isotopologs and highly excited vibrational satellites. The frequency ranges covered were $440\text{--}540$, $619\text{--}631$, $638\text{--}648$, $770\text{--}855$, $875\text{--}930$, $967\text{--}1050$, $1083\text{--}1093$, $1100\text{--}1159$, $1168\text{--}1198$, $1576\text{--}1591$,

and $1614\text{--}1627 \text{ GHz}$, as in our previous analysis [27]. In addition, we obtained new data at $530\text{--}595 \text{ GHz}$ and data with improved S/N in the $780\text{--}835 \text{ GHz}$ region. Most of these multiplier sources were previously described [58]. However, the multiplier chain with frequency range coverage between $967\text{--}1050 \text{ GHz}$ was not described in that work. This chain consists of two cascaded triplers after the amplified W-band stage, the peak output power is near $100 \mu\text{W}$. The efficiency of frequency multipliers is usually strongly frequency-dependent. In addition, recording conditions and sensitivities of detectors can have strong influences on the quality of the spectra. Particularly good S/N ratios were reached around 600 , 800 , 900 and $1100\text{--}1200 \text{ GHz}$. The S/N ratios changed considerably within each scan and were usually lower towards the edges. Uncertainties of 50 kHz were assigned to isolated lines with good to moderate S/N ratios, smaller uncertainties, down to 10 kHz , in cases of very good S/N and very symmetric lineshapes, larger uncertainties if the S/N was poorer or if the lineshape was rather asymmetric, e.g. because of the proximity to another line.

2.2. Rotational spectra at the Universität zu Köln

All measurements at the Universität zu Köln were recorded at room temperature in static mode employing different Pyrex glass cells having an inner diameter of 100 mm with pressures in the range of $0.1\text{--}0.5 \text{ Pa}$ at lower frequencies, increased to around 1.0 Pa at higher frequencies. The window material was Teflon at lower frequencies, whereas high-density polyethylene was used at higher frequencies. Frequency modulation was used throughout with demodulation at $2f$, causing an isolated line to appear close to a second derivative of a Gaussian.

The $J = 2 - 1$ and $3 - 2$ transitions of several vibrational states around 37 and 55 GHz , respectively, were recorded with an Agilent E8257D microwave synthesizer as source and a home-built Schottky diode detector. A 7 m long double pass absorption cell was used for these measurements. All other measurements were carried out in shorter (around 3 m) single pass cells. The $J = 4 - 3$ transitions of $v_8 = 2$ around 74.0 GHz were recorded using a backward-wave oscillator (BWO) based 4 mm synthesizer AM-MSP 1 (Analytik & Meßtechnik GmbH, Chemnitz, Germany) as source and a Schottky-diode as detector. A small number of methyl cyanide rotational transitions were recorded around 850 GHz and around 880 GHz with the Cologne Terahertz Spectrometer [59] using a BWO as source and a liquid helium cooled InSb hot-electron bolometer (QMC) as detector. Further, extensive measurements were carried out in parts of the $1330\text{--}1501 \text{ GHz}$ region employing a frequency multiplier (Virginia Diode, Inc.) driven by an Agilent E8257D microwave synthesizer as source and the InSb bolometer as detector. This spectrometer was used previously for investigations of CH_3SH and $^{13}\text{CH}_3\text{OH}$ [60, 61].

2.3. Infrared spectrum

One infrared spectrum of CH_3CN was recorded between 600 and 989 cm^{-1} at 0.0016 cm^{-1} resolution using an MCT detector with a Bruker 120 HR Fourier transform spectrometer at Pacific Northwest National Laboratory (PNNL); the $2\nu_8$ region is

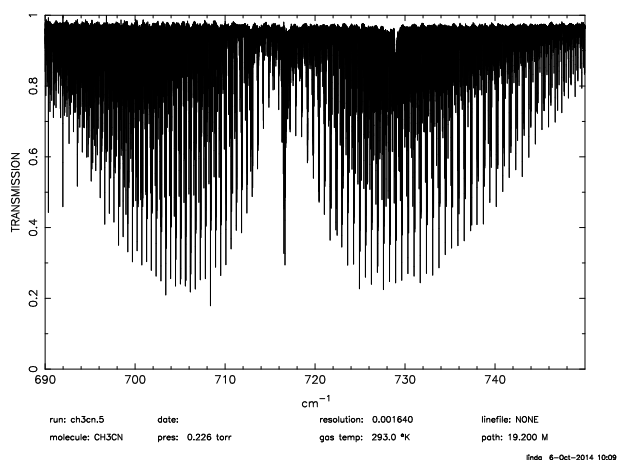


Figure 1: The $2\nu_8$ spectrum of CH_3CN recorded at 0.0016 cm^{-1} resolution using a Bruker 120 HR FTS at PNNL; the optical path is 19.2 m, and the gas pressure is 0.226 Torr at 293 K. The sample included $\sim 0.4\%$ CO_2 and $\sim 3\%$ OCS for frequency calibration.

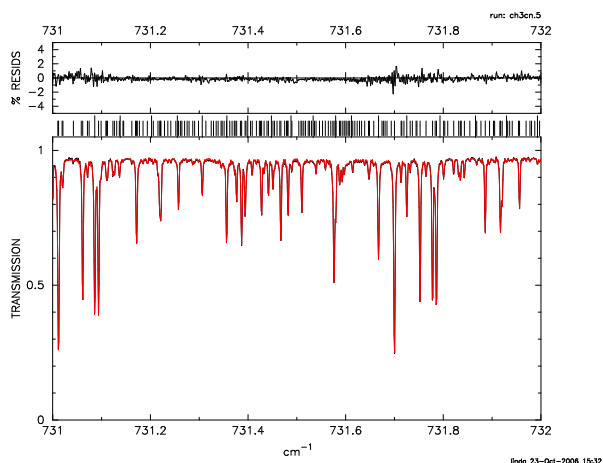


Figure 2: Comparison of CH_3CN lines at 731.5 cm^{-1} retrieved by least squares curve-fitting. The lower panel has the observed spectrum in black overlaid by the synthetic spectrum in red. The upper panel shows the residual difference (calc.-obs.) in percent. The optical path is 19.2 m, and the gas pressure is 0.226 Torr at 293 K. The vertical lines between the two panels mark the measured features (every 10th line is taller).

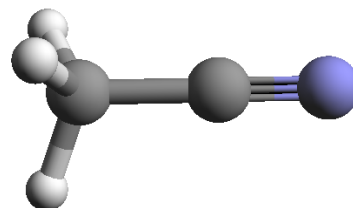


Figure 3: Model of the methyl cyanide molecule. The C atoms are shown in gray, the H atoms in light gray, and the N atom is shown in blue. The a -axis is along the CCN atoms and is also the symmetry axis.

shown in **Fig. 1**. For this, a multi-pass absorption cell, set to an optical path length of 19.20 m, was filled to 0.226 Torr (30.1 Pa) of CH_3CN at 293.0 K. Small amounts of OCS and CO_2 were added to the sample for frequency calibration; comparison of 63 well-isolated lines of the ν_3 fundamental of OCS at 860 cm^{-1} [62] produced a calibration factor 1.000000868 (21) with an rms of 0.0000182 cm^{-1} . Nearly 12000 line positions and relative intensities between 690 and 765 cm^{-1} were retrieved using non-linear least-squares curve-fitting [63]. However, the corresponding precisions in the $2\nu_8$ region were an order of magnitude worse than the OCS calibration because the observed spectrum was so congested. **Fig. 2** shows the quality of the curve fitting for a small section of $2\nu_8$.

Two other room-temperature spectra spanned the ν_8 region (250 to 400 cm^{-1}) at 294 K using pressures of 0.78 Torr (103.9 Pa) with a 19.2 m path and 1.01 Torr (134.5 Pa) with a 3.24 m path. These were calibrated using residual water inside the FTS [64] to obtain a calibration factor of 1.000000207 (495) with an rms of only 0.000140 cm^{-1} . The spectra were very congested and no attempt was made to retrieve positions and intensities by curve-fitting. Previously reported CH_3CN line centers [33] had better precisions in this second region, and so these new ν_8 data were used mainly to search for hot bands and to confirm assignments in the $2\nu_8$ region.

3. Results

3.1. Overview of CH_3CN spectroscopy

As shown in **Fig. 3**, the light H atoms are the only ones not on the symmetry axis of CH_3CN . Hence, methyl cyanide is a strongly prolate molecule ($A \gg B$). Rotational transitions with $\Delta K = 0$ are strong because of the large dipole moment of $3.92197(13)\text{ D}$ [65]. $\Delta K = 3$ transitions only gain intensity through centrifugal distortion effects and are usually very weak. In the absence of perturbations, the purely axial parameters A (or $A - B$), D_K , etc. cannot be determined by rotational spectroscopy. Moreover, even rovibrational spectroscopy yields, strictly speaking, the differences ΔA (or $\Delta(A - B)$), ΔD_K , etc. from single state analyses, but not the ground state parameters. In the case of CH_3CN , analyses of three IR bands involving two doubly degenerate vibrational modes, ν_8 , $\nu_7 + \nu_8$, and $\nu_7 + \nu_8 - \nu_8$, permitted the axial parameters to be determined [32].

Table 1: Energies (cm^{-1}) and symmetries Sym of low-lying vibrational states of methyl cyanide.

State	Sym.	Energy	Reference
$v = 0$	A	0.0	per definitionem
$v_8 = 1$	E	365.024	this work
$v_8 = 2^0$	A	716.750	this work
$v_8 = 2^2$	E	739.148	this work
$v_4 = 1$	A	920.290	Ref. [45]
$v_7 = 1$	E	1041.855	Ref. [45]
$v_8 = 3^1$	E	1077.786	Ref. [45]
$v_8 = 3^3$	A	~ 1122.15	Ref. [45]
$v_4 = v_8 = 1$	E	1290.0	Ref. [69]
$v_3 = 1$	A	1385.2	Ref. [69]
$v_7 = v_8 = 1^0$	A	1401.7	Ref. [69]
$v_7 = v_8 = 1^2$	E	1408.9	Ref. [69]
$v_8 = 4^0$	A	$\sim 1426.$	this work
$v_8 = 4^2$	E	1447.9	Ref. [69]
$v_6 = 1$	E	1449.7	Ref. [69]
$v_8 = 4^4$	E	$\sim 1514.$	this work

Methyl cyanide has four totally symmetric and four doubly degenerate vibrational modes. The three lowest energy fundamentals are the doubly degenerate CCN bending mode ν_8 at 365.024 cm^{-1} [33], the totally symmetric CC stretching mode ν_4 at 920.290 cm^{-1} [45], and the doubly degenerate CH_3 rocking mode ν_7 at 1041.855 cm^{-1} [45]. They are comparatively weak with integrated room temperature cross sections in the range 2 to $5 \times 10^{-19} \text{ cm}^2/\text{molecule}$ [66, 67]. Interestingly, the $2\nu_8$ overtone band is slightly stronger than ν_8 or ν_4 [66, 67]. In the case of ν_4 , line-by-line intensity measurements have also been published [68]. A list of low-lying vibrational states, up to $\sim 1500 \text{ cm}^{-1}$, is given in **Table 1**.

The fairly large value of A , $\sim 5.27 \text{ cm}^{-1}$, leads to a rapid increase in rotational energy with K which is shown up to $\sim 1700 \text{ cm}^{-1}$ in **Fig. 4**. The highest K levels observed involve $K = 21$, and the $J = K = 21$ energy is at 2313 cm^{-1} .

The vibrational energy of the lowest, doubly degenerate vibrational mode $\nu_8 = 1$ is at 365.024 cm^{-1} [33], similar to the $J = K = 8$ energy in $v = 0$ (**Fig. 4**). The Coriolis interaction between the two l components of $\nu_8 = 1$, described by $\zeta \approx 0.877$, is large, close to the limiting case $\zeta = 1$. It effectively shifts the levels of the $l = +1$ ladder to lower energies and the energies of the $l = -1$ ladder to higher energies. As a result, $K = 1$ of $l = +1$ is actually lower than $K = 0$. More importantly, levels with $\Delta K = \Delta l = 2$ are near-degenerate, leading to pronounced q_{22} interaction within $\nu_8 = 1$. The interaction is highlighted in **Fig. 4** for $K = 4$ of $\nu_8 = 1^{+1}$ and $K = 2$ of $\nu_8 = 1^{-1}$. As a consequence of the Coriolis interaction within $\nu_8 = 1$, the K level spacing in $\nu_8 = 1^{+1}$ is considerably smaller than in $v = 0$ while it is larger in $\nu_8 = 1^{-1}$. Šimečková et al. [41] found evidence of a resonant $\Delta\nu_8 = \pm 1$, $\Delta K = \mp 2$, $\Delta l = \pm 1$ interaction between $v = 0$, $K = 14$ and $\nu_8 = 1^{+1}$, $K = 12$ (**Fig. 4**) at $J = 43$. They analyzed the interaction approximately because only two slightly perturbed transitions in the ground vibrational state ($J'' = 38$ and 39) were observed.

Other global interactions within an isolated degenerate vibrational mode of a strongly prolate rotor are rare, in contrast

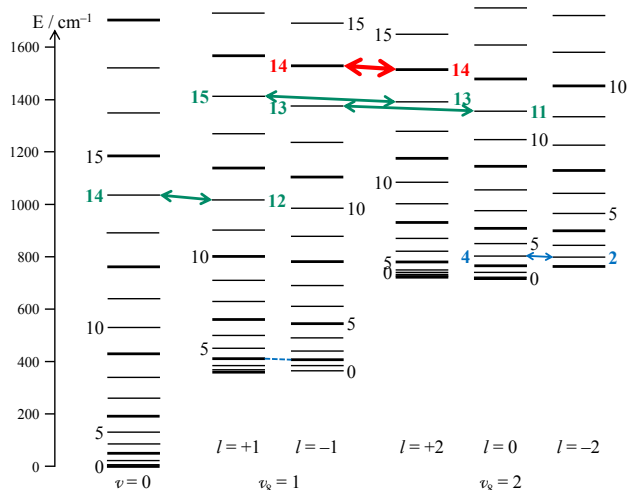


Figure 4: K level structure of methyl cyanide displaying the $J = K$ energies for $\nu_8 = 0, 1$, and 2 with their l substates. Levels with A symmetry are drawn with thick lines. Coriolis coupling within $\nu_8 = 1$ causes near-degeneracies of levels having $\Delta K = \Delta l = 2$ and hence strong q_{22} interaction shown exemplarily for one pair of low K by a thin, dashed (blue) line. The same applies to $\nu_8 = 2^{\pm 2}$, where a q_{22} resonance occurs between $K = 4$ and $K = 2$ of the $l = 0$ and -2 components, respectively (thin blue left right arrow). The spacing between K levels is fairly large compared to the vibrational spacing causing a $\Delta K = 0$, $\Delta l = 3$ Fermi interaction at $K = 14$ between $\nu_8 = 1^{-1}$ and $\nu_8 = 2^{+2}$, displayed by a thick (red) left right arrow. Also shown are $\Delta K = \mp 2$, $\Delta l = \pm 1$ interactions between $\nu_8 = 0$ and 1 as well as between $\nu_8 = 1$ and 2 (medium thick green left right arrow).

to nearly spherical or oblate rotors, see, e.g. Ref. [70] and references therein.

The basic features of the K level structure of $\nu_8 = 2$ are similar to those of $\nu_8 = 1$, the main difference being three l components, $l = 0, \pm 2$ instead of two. Pronounced q_{22} interaction occurs also within $\nu_8 = 2$, in particular between $l = 0$ and $l = -2$ at low K with a resonant interaction $K = 4$ of $\nu_8 = 2^0$ and $K = 2$ of $\nu_8 = 2^{-2}$, as shown in **Fig. 4**. In addition, the K level structure in $\nu_8 = 2^{+2}$ is irregular at low K with $K = 2$ being lowest in energy and $K = 3$ still slightly below $K = 0$. Such irregular K level structure in states such as $\nu_8 \geq 2$ may lead to *local* perturbations within the vibrational state with, e.g., $\Delta K = \pm 1$, $\Delta l = \mp 2$ (q_{21}), $\Delta K = \pm 3$, $\Delta l = 0$ (ϵ), etc. In the case of several different such local perturbations, indeterminacies, as investigated in Ref. [70] and references therein, may need to be considered. A global effect may be caused by a $\Delta K = \pm 4$, $\Delta l = \pm 4$ interaction (f_{44}). However, this effect is rather small in a strongly prolate rotor, and f_{44} was barely determined for $\nu_{10} = 2$ of propyne [51].

3.2. $v = 0$

Recently, we have revisited the ground state rotational spectrum of several isotopic species of methyl cyanide [27]. Besides new data, ground state transition frequencies for the main species were mainly taken from Ref. [28] with additional lower frequency data [29, 30, 31]. The J and K quantum numbers extended to 89 and 21, respectively, see also **Table 2**.

Table 2: Maximum J and K values and number of rotational and IR lines in vibrational substates of methyl cyanide used in the present / previous study^a, with numbers retained in the present fit from previous data in parentheses and rms error for subsets of data.

	$v = 0$	$v_8 = 1^{+1}$	$v_8 = 1^{-1}$	$v_8 = 2^0$	$v_8 = 2^{+2}$	$v_8 = 2^{-2}$
$J_{\max}(\text{rot})$	89/89 (89)	88/75 (45)	88/75 (8)	88/68 (20)	88/68 (20)	88/68 (20)
$K_{\max}(\text{rot})$	21/21 (21)	19/13 (7)	17/11 (7)	19/12 (12)	13/8 (7)	11/8 (6)
no. of rot lines ^b	316/225 (196)	596/c (58)	527/c (37)	424/85+5 ^d (18)	546/73+6 ^d (24)	299/45+5 ^d (9)
rms error ^e	0.859/0.911	0.807/-	0.796/-	0.982/-	0.915/-	0.957/-
$J_{\max}(\text{IR})$	-/- (-) ^f	71/71 (71)	74/74 (74)	66/65	53/37	66/57
$K_{\max}(\text{IR})$	7/7 (7)	11/11 (11)	11/11 (11)	12/10	2/4	7/4
no. of IR lines ^b	5/5 (5) ^f	836/g (836)	861/g (861)	935/944 (0)	40/44 ^h (0)	197/122 ^h (0)
rms error ^e	0.716/0.643	0.826/-	0.797/-	0.771/-	0.849/-	0.727/-

^a Previous data from Refs. [27, 33, 42] and references therein for $v = 0$, $v_8 = 1$, and $v_8 = 2$, respectively.

^b Each blend of lines counted as one line. Hyperfine splitting in the present and previous data was considered in the present fit. Transitions with $K = 0$ of l -doubled states associated with $v_8 = 1^{-1}$ and $v_8 = 2^{+2}$, respectively, in the present data set. Lines, which were weighted out, are not counted. Transitions between vibrational states were counted for higher vibrational state.

^c Only total number given as 187.

^d Numbers after the plus were given zero weight in Ref. [42].

^e No rms error was given in Refs. [33, 42]; we assume parameter uncertainties are based on standard errors, i.e. the rms error is 1.0 by definition. The rms errors for data from Ref. [27] were calculated from the experimental data available at <http://www.astro.uni-koeln.de/site/vorhersagen/daten/CH3CN/CH3CN/gs/>.

^f No individual $\Delta K = 3$ loops were given in Ref. [32]. We used the five $\Delta K = 3$ splittings from Table II in that work with the reported uncertainties.

^g Only total number given as 1705. Slight difference to our numbers mainly because of blended lines, see subsection 3.3.2.

^h 52 assigned $2\nu_8 - \nu_8$ hot band transitions are mentioned in Ref. [42]; see subsection 3.4.3.

The first objective and the main achievement of the present study with respect to the ground vibrational state was the detailed analysis of the resonance between $v = 0$ and $v_8 = 1$ mentioned in the previous subsection. Absorption features pertaining to $v = 0$, $K = 14$, $J'' = 42$ to 49 and to $v_8 = 1^{+1}$, $K = 12$, $J'' = 41$ to 45 and 47 to 49 could easily be assigned in our spectral recordings because of their intensities. Moreover, the perturbed transitions showed regular and mostly small displacements from predictions which did not take the interaction into account. The largest shifts were ~ 24 MHz and ~ 16 MHz for the $J = 43 - 42$ and $44 - 43$ transitions, respectively. After these data had been included in the fit, two pairs of cross ladder transitions between the two vibrational states were predicted to have sufficient intensities for observation. The weaker pair with $J = 44 - 43$ were identified, but the ones with $J = 43 - 42$, predicted to be stronger by about a factor of 2.5 than the former ones, were blended with stronger lines. **Fig. 5** shows a detail of the energy levels around the levels perturbed most by the interaction between $v = 0$ and $v_8 = 1$ as well as the observed transitions.

The measurements in Köln between 1.33 and 1.5 THz closed a considerable part of the gap between ~ 1.2 and ~ 1.6 THz in the previous data set [27]. We reanalyzed also the spectra between ~ 1.1 and 1.2 THz taken at JPL and realized that the uncertainties were very similar to those of Ref. [28], albeit with a slightly better coverage and extending to slightly higher K . Therefore, we employed the JPL data in our present analysis.

We also attempted to assign rotational transitions with $\Delta K = 3$. Their predicted uncertainties were small enough to identify them easily if the distortion dipole moment was as large as 10^{-4} D. However, even then they could be blended with stronger methyl cyanide transitions, of which several are not

yet assigned. No assignments could be made, leading us to conclude that the distortion dipole moment is probably smaller. It is also possible that the calculated uncertainties of these transitions are too small.

3.3. $v_8 = 1$

3.3.1. Rotational data

The rotational spectrum of CH_3CN in its $v_8 = 1$ lowest vibrational state ($E_{\text{vib}} = 365.02 \text{ cm}^{-1}$) was reported on several occasions [35, 36, 37, 38]. The $J'' = 19$ data [37] had residuals in part much larger than the reported 50–150 kHz uncertainties already in the initial report, which were confirmed here. Therefore, and because of their more limited K range, we omitted these data from our line list. The laser side-band measurements [38] were also quite limited in K and in accuracy (around 1 MHz) and were superseded by our more accurate data. Part of the lower frequency data up to 150 GHz [35, 36] was also superseded by our data, the remainder was retained in the fit with reported uncertainties. We also used in our fit direct l -type transitions between the $k \times l = +1$ components [39, 40].

The interaction between $v = 0$ and $v_8 = 1$ was already described in subsection 3.2. Interactions of $v_8 = 1$ with any higher vibrational states have not been reported in the literature prior to the present investigation.

Based on the studies on propyne [50, 51], we expected an interaction between $v_8 = 1^{-1}$ and $v_8 = 2^{+2}$ at K slightly higher than 12 because of the slightly larger vibrational energy difference in methyl cyanide compared to propyne, 374.12 cm^{-1} compared to 340.41 cm^{-1} . Inspection of the energy levels of these CH_3CN vibrational states indicated a resonant interaction to occur at $K = 14$, see **Figs. 4 and 6**. Using the value of the main Fermi parameter from the propyne study [50, 51] was

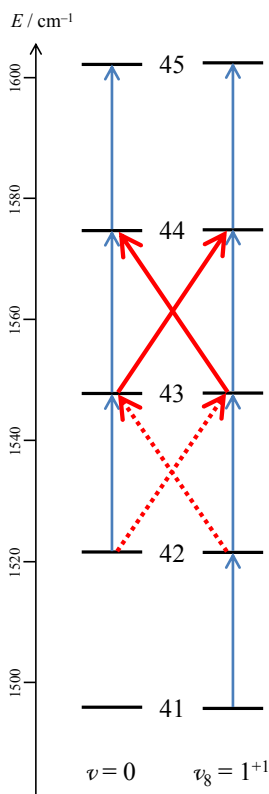


Figure 5: Detail of the energy level structure of methyl cyanide around the perturbation between the ground state and $v_8 = 1$. Transitions observable in our spectral recordings are represented by arrows; those with solid lines indicate transitions in the final line list, including two transitions between the vibrational states. Two further transitions between $v = 0$ and $v_8 = 1$ with significant intensities, but overlapped with stronger transitions are displayed by arrows with dotted lines.

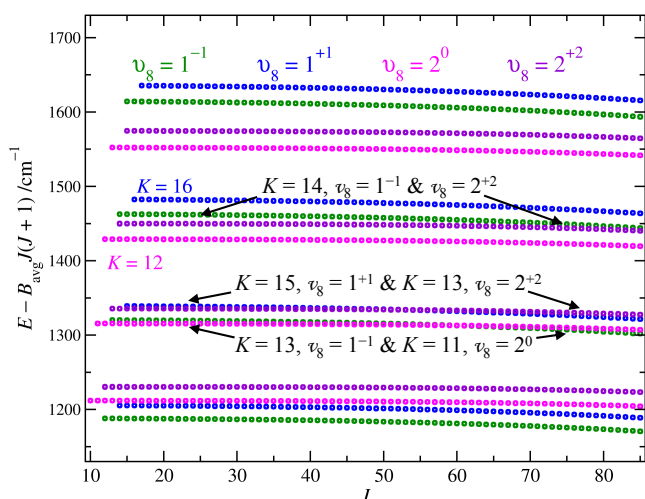


Figure 6: Reduced energy plot in the region of the perturbations between $v_8 = 1$ and 2. A $\Delta v_8 = \pm 1$, $\Delta K = 0$, $\Delta l = \pm 3$ Fermi resonance occurs near 1450 cm^{-1} . In addition, two resonant $\Delta v_8 = \pm 1$, $\Delta K = \mp 2$, $\Delta l = \pm 1$ interactions are indicated near 1330 cm^{-1} .

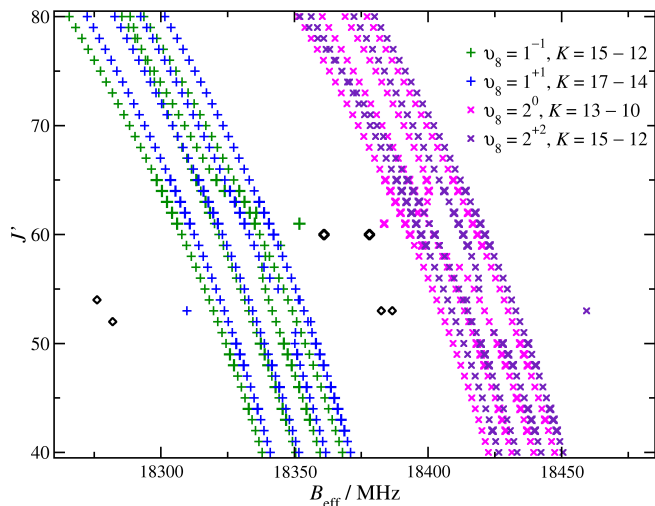


Figure 7: Fortrat diagram in the region of the resonant $\Delta v_8 = \pm 1$, $\Delta K = \mp 2$, $\Delta l = \pm 1$ interactions between $v_8 = 1$ and 2. Observed transitions are indicated by larger symbols.

a reasonable initial guess along with published parameters for $v_8 = 2$. Our final results suggest the interaction to be strongest at $J = 91$ and 92 , however, the highest assignments presently reach $J = 65 - 64$ for either state at these rather high values of K . The magnitudes of the perturbations between $v_8 = 1^{-1}$ and $v_8 = 2^{+2}$ showed strong variations with K , as expected, but also with J , most likely because the energy differences between the two vibrational ladders also change significantly with J . The perturbations exceed 1 MHz at $J = 24 - 23$, $K = 10$ and reach 37 MHz at $K = 14$. At $J = 65 - 64$, perturbations in the spectrum larger than 1 MHz were already seen at $K = 5$, and the shift exceeded 300 MHz at $K = 14$.

The data set up to $J = 28 - 27$ revealed perturbations in $K = 15$ of $v_8 = 1^{+1}$ which were not accounted for by this $\Delta v_8 = \pm 1$, $\Delta K = 0$, $\Delta l = \pm 3$ Fermi interaction. They were caused by a $\Delta v_8 = \pm 1$, $\Delta K = \mp 2$, $\Delta l = \pm 1$ interaction with $K = 13$ of $v_8 = 2^{+2}$ (Figs. 4 and 6), the same type of interaction as between $v = 0$ and $v_8 = 1$. It is shown in Fig. 7 that the strongest mixing occurs at $J = 52$ and 53 . Among the four transitions with $J = 53 - 52$, the $v_8 = 1$ transition was too close to the lower frequency limit of its experimental spectral window, the $v_8 = 2$ transition was observed with a total perturbation of 1.8 GHz . In addition, both cross ladder transitions were observed, but one was too close to a stronger line to determine a reliable and sufficiently accurate transition frequency. In addition, the $J = 54 - 53$ transition frequencies were determined for both vibrational states.

The same type of interaction showed noticeable effects in the spectrum at higher J (up to 49) between $K = 13$ of $v_8 = 1^{-1}$ and $K = 11$ of $v_8 = 2^0$, see also Figs. 4 and 6, with the largest perturbation at $J = 60$, as shown in Fig. 7. The regular rotational transitions were detected for both vibrational states in the case of $J = 61 - 60$, both cross-ladder transitions were detected for $J = 60 - 59$. The total perturbations amounted to more than 1.5 GHz for the former two transitions and more than 3.0 GHz

for the latter two. The remaining associated transitions were about an order of magnitude weaker and were not identified.

We mention for completeness a $\Delta v_8 = \pm 2$, $\Delta K = \mp 2$, $\Delta l = \pm 4$ interaction between $K = 15$ of $v_8 = 1^{-1}$ and $K = 13$ of $v_8 = 3^{+3}$ whose strongest effects are currently predicted at $J = 73/74$. Should the interaction cause any measurable perturbation in the spectrum, it should be very localized because of the high order of the interaction. A $\Delta v_8 = \pm 2$, $\Delta K = \mp 1$, $\Delta l = \pm 2$ interaction is possible between $K = 26$ of $v_8 = 1^{-1}$ and $K = 25$ of $v_8 = 3^{+1}$, but these levels are so high in energy that they are unlikely to be accessed.

The overall coverage of our $v_8 = 1$ is indicated in **Table 2**. Besides the few cross-ladder transitions between $v_8 = 1$ and 2, the data are limited almost entirely to regular rotational transitions. Transitions between different l components may occur because of $\Delta l = \Delta K = 2$ mixing facilitated by q . Such transitions were predicted to be very weak, and only one transition was identified, so we view this identification as tentative. We also searched for $\Delta K = 3$ transitions, but, as in the case of the ground vibrational state, were unable to identify any with sufficient certainty.

3.3.2. The v_8 band

An investigation of v_8 was reported in Ref. [33], and we use the line list from that work as indicated in subsection 2.3. Lines with full weight were assigned uncertainties of 0.0002 cm^{-1} , except for the small number of ${}^rQ_{10}$ lines, for which 0.0004 cm^{-1} were used because of their persistent larger residuals in the fits. Lines with weight 0.1 were assigned uncertainties of 0.0006 cm^{-1} . A few lines (around 10 to 20) with weight 0.1 were omitted either because of very large residuals or because the line was blended, but was the weaker line, and the stronger line was not in the line list. In other rare cases, blended lines (12×2) were treated as intensity-weighted averages. The effects caused by these differences in the previous and present line list on the parameter values, uncertainties and on the quality of the fit are probably negligible.

3.4. $v_8 = 2$

3.4.1. Rotational data

The rotational spectrum of methyl cyanide in its $v_8 = 2$ excited vibrational state was also studied extensively. As for $v_8 = 1$, the lower frequency data up to 112 GHz [36, 46, 47] were in part superseded by new data taken in Cologne; see **Fig. 8** for an overview of the $J = 2 - 1$ transition. The remaining data were used in the final fit with reported uncertainties. The $J'' = 19$ data from Ref. [37] had reported uncertainties of 50–150 kHz, though not specified. They fit on average to within 100 kHz and were retained in the final fit with assigned uncertainties of 100 kHz. The laser-sideband data [38] were not retained in the fit for the same reasons as in the case of $v_8 = 1$.

Interactions between $v_8 = 1$ and 2 were described in subsection 3.3.

The energy difference between $v_8 = 2^0$ and $v_4 = 1$ is 203.54 cm^{-1} , rather large for a resonant Fermi interaction to occur. A $\Delta K = \pm 1$, $\Delta l = \mp 2$ interaction between $K = 6$ of

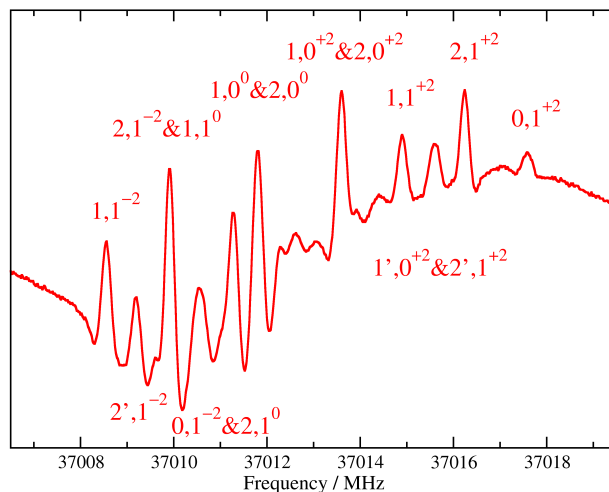


Figure 8: The rotational spectrum of CH_3CN in the region of $v_8 = 2$, $J = 2 - 1$ with largely resolved ^{14}N hyperfine splitting. All but one feature (near 37010.5 MHz) used in the fit are indicated by assignment labels centered on the respective line. The first number gives the lower state F value. Primed F values indicate the weaker $\Delta F = 0$ transitions; unprimed have $\Delta F = 1$. Numbers after the comma signal the K value and, as superscripts, the l value. An ampersand was used to indicate overlapping HFS components.

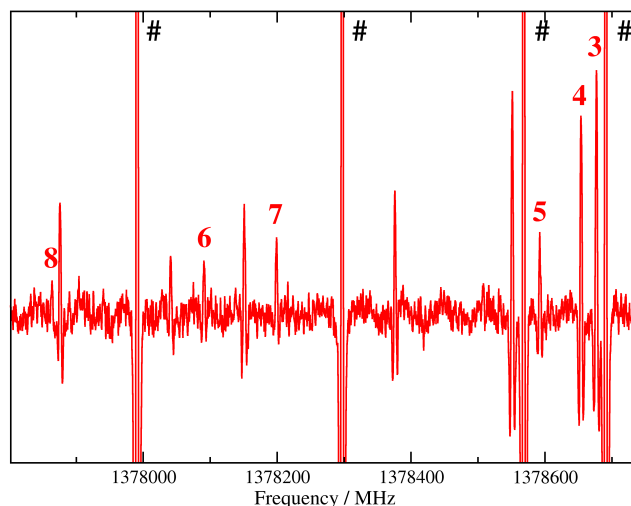


Figure 9: The rotational spectrum of CH_3CN in the region of $J = 75 - 74$ displaying the perturbation of $v_8 = 2^{-2}$, $K = 6$ by $v_4 = 1$, $K = 5$. The K values are given for $v_8 = 2^{-2}$; the perturbation shifts the $K = 6$ transition lower by 350 MHz. The clipped lines marked with a black pound sign are (from left to right) $K = 4, 3, 2$ and 1 transitions of $v_8 = 1^{+1}$.

$\nu_8 = 2^{-2}$ and $K = 5$ of $\nu_4 = 1$ was considered in the analysis of the ν_4 band [45]. We determined a resonance to occur at $J = 73$. Our closest $\nu_8 = 2$ transition is $J = 75 - 74$ which is calculated to be shifted by ~ 350 MHz, see **Fig. 9**. A $\Delta K = \pm 3$ ($K = 8/5$), $\Delta l = 0$ interaction, resonant at $J = 57$, was identified here for the first time. Our closest $\nu_8 = 2$ transitions are $J = 56 - 55$ and $J = 59 - 58$ with calculated shifts of about 20 MHz and 25 MHz, respectively.

Even though the energy difference between $\nu_8 = 2^2$ and $\nu_7 = 1$ is 302.71 cm^{-1} , much larger than the difference between $\nu_8 = 2^0$ and $\nu_4 = 1$, the Coriolis interaction *within* the l components of $\nu_8 = 2^2$ and $\nu_7 = 1$ lead to a $\Delta K = 0$, $\Delta l = 3$ Fermi resonance at $K = 13$ of $\nu_8 = 2^{-2}$ and $\nu_7 = 1^{+1}$, which caused perturbations in the ν_7 band [45]. We were able to assign transitions up to $K = 11$ in $\nu_8 = 2^{-2}$. Effects of this Fermi resonance were clearly present, but they were not sufficient to separate them reliably from the stronger Fermi resonance of $\nu_8 = 2^{-2}$ with $\nu_8 = 3^{+1}$ occurring at $K = 12$.

A $\Delta K = \pm 2$, $\Delta l = \mp 1$ resonant interaction between $K = 9$ of $\nu_8 = 2^{-2}$ and $K = 7$ of $\nu_7 = 1^{-1}$ near $J = 88$ shows no effects in the spectrum up to $J = 65$. At higher J , the intensities of these high energy lines are just too small to be identified reliably in our present spectral recordings. An identical type of interaction may occur at $K = 14$ and 15 of $\nu_8 = 2^0$ and $K = 12$ and 13 of $\nu_7 = 1^{+1}$, however, we have no assignments yet for such high energy transitions.

Fermi interactions between $\nu_8 = 1$ and 2 may also occur between $\nu_8 = 2$ and 3 . There are $\Delta \nu_8 = \pm 1$, $\Delta K = 0$, $\Delta l = \pm 3$ Fermi resonances between $\nu_8 = 2^{-2}$ and $\nu_8 = 3^{+1}$ at $K = 13$ and between $\nu_8 = 2^0$ and $\nu_8 = 3^{+3}$ at $K = 14/15$; a resonant interaction may occur for the latter at $K = 15$, currently predicted at $J = 71/72$. In the case of $\nu_8 = 2^{-2}$, our assignments extend to $K = 11$ and to $K = 13$ for $\nu_8 = 2^0$. These levels are perturbed sufficiently, up to more than 20 MHz and almost 30 MHz, respectively, to evaluate the corresponding Fermi parameters.

$\Delta \nu_8 = \pm 1$, $\Delta K = \mp 2$, $\Delta l = \pm 1$ interactions are expected to occur between $\nu_8 = 2$ and 3 at $K = 12$, $l = -2$ and $K = 10$, $l = -1$; at $K = 14$, $l = 0$ and $K = 12$, $l = +1$; and at $K = 16$, $l = +2$ and $K = 14$, $l = +3$. One tentative assignment exists for the latter interaction. In addition, we have 6 assignments extending to as high as $K = 19$ of $\nu_8 = 2^{+2}$, but several have residuals slightly larger than the estimated uncertainties, which may be explained by unregonized overlap with other transitions, by perturbations of one or both rotational levels, or by underestimation of the uncertainties.

We should point out that there may be interactions even higher in ΔK or Δl , but we do not have evidence or data for these at present. No attempts were made to locate interactions of $\nu_8 = 2$ with vibrational states higher than $\nu_8 = 3$. Such interactions may be non-negligible, but most likely at K levels with few or no assignments in our present data set.

3.4.2. The $2\nu_8$ band

Transition frequencies were determined entirely from our new IR spectrum since Ref. [42] did not report their experimental values.

The strongly allowed transitions in the $2\nu_8$ IR band are those with $\Delta K = \Delta l = 0$. The K quantum numbers reach 12, higher by 2 than in the previous study [42], while the maximum value of J is essentially identical, 66 versus 65. As in the previous study, we were unable to make any assignments in the perpendicular $\Delta l = \pm 2$, $\Delta K = \mp 1$ component, which is only weakly allowed, but we were able to make numerous assignments to levels in the $l = \pm 2$ component which gain intensity through $\Delta K = \Delta l = 2$ mixing. This mixing is particularly strong for $K = 4$ of $\nu_8 = 2^0$ and $K = 2$ of $\nu_8 = 2^{-2}$ with the former being lower in energy at low- J until the energy ordering changes at $J = 40/41$. As a consequence, many ${}^oR_4(J)$ and ${}^oP_4(J)$ IR transitions were assigned covering most of the J from 8 to 62. It is interesting to note that the mixing of these K levels changes only slightly between adjacent J levels such that rotational transitions between these mixed K levels are too weak to be observed with the exception of those at the change of the energy ordering. As also shown in **Table 2**, the IR assignments to $\nu_8 = 2^{-2}$ cover mostly $K = 1$ to 3 with several $K = 4$ assignments and few higher ones, those to $\nu_8 = 2^{+2}$ cover mostly $K = 0$ to 2 ; a few sparse higher- K assignments are viewed as tentative and were weighted out in the final fit.

Isolated lines with sufficiently good signal-to-noise ratio were given uncertainties of 0.0002 cm^{-1} , as in the case of ν_8 . Noisier lines or lines in more crowded regions, such as in the Q -branch, were given larger uncertainties of 0.0003 cm^{-1} to 0.0010 cm^{-1} .

We were also able to make assignments in the $3\nu_8 - \nu_8$ hot band and in the $\nu_7 - \nu_8$ difference band, as was done earlier [42]. The latter band gains intensity because of an anharmonic resonance between $\nu_7 = 1$ and $\nu_8 = 3$ [42]. We will publish our detailed account on these states separately. Judging from the observation of transitions in $2\nu_8^0$ with $K = 12$, we may be even able to make assignments in $4\nu_8 - 2\nu_8$.

3.4.3. The $2\nu_8 - \nu_8$ hot band

Because of the strength of the $2\nu_8^0$ band, little can be gained from the analysis of the $2\nu_8^0 - \nu_8$ hot band. On the other hand, IR transitions to $2\nu_8^2$ are limited to those which gain intensity through q_{22} interaction, see previous subsection. These transitions belong to the $l = -2$ component for the most part. According to Koivusaari et al. [42], transitions of $2\nu_8^2 - \nu_8$ are frequently overlapping each other or are blended with cold band transitions. Nevertheless, they made 52 r -branch assignments with $3 \leq K'' \leq 7$. 50 of these assignments were made in the $l = +2$ component of $2\nu_8^2 - \nu_8$ of which 6 had zero weight. Since no IR transition frequencies were published in Ref. [42], we searched our ν_8 spectra for suitable hot band transitions. Unfortunately, the congestion of the spectral lines and the quality of the spectra were such that only few lines could be assigned with certainty.

3.5. Remarks on higher vibrational states

The vibrational state $\nu_8 = 2$ interacts with $\nu_4 = 1$, with $\nu_7 = 1$, and with $\nu_8 = 3$. Therefore, we need to take into account previously determined parameters for these states [45], which involve interactions among these three vibrational states.

Table 3: Spectroscopic parameters or differences Δ thereof ^{a,b} (cm⁻¹, MHz)^c of methyl cyanide within vibrational states $v_8 \leq 2$.

Parameter X	$v = 0$	$v_8 = 1$	$v_8 = 2^0$	$v_8 = 2^2$
E_{vib}^c	0.0	365.024 365 (9)	716.750 42 (13)	739.148 225 (56)
$(\Delta)^b (A - B)$	148 900.103 (66)	-115.930 (26)	-187.404 (18)	-259.956 (122)
$(\Delta)^b B$	9 198.899 167 (11)	27.530 277 (49)	54.057 316 (111)	54.502 729 (70)
$(\Delta)^b D_K \times 10^3$	2 830.6 (18)	-11.46 (48)	-20.19 (34)	-7.45 (164)
$(\Delta)^b D_{JK} \times 10^3$	177.407 87 (25)	0.987 48 (60)	1.675 5 (25)	1.808 8 (15)
$(\Delta)^b D_J \times 10^6$	3 807.576 (8)	95.599 (17)	216.319 (37)	189.162 (31)
$(\Delta)^b H_K \times 10^6$	164.6 (66)	14.9 (22)		
$(\Delta)^b H_{KJ} \times 10^6$	6.062 0 (14)	0.034 1 (22)	0.150 3 (239)	0.024 9 (55)
$(\Delta)^b H_{JK} \times 10^9$	1 025.69 (15)	2.59 (6)	17.71 (34)	-0.37 (21)
$(\Delta)^b H_J \times 10^{12}$	-237.4 (21)	315.3 (30)	200.7 (63)	627.8 (59)
$(\Delta)^b L_{KKJ} \times 10^{12}$	-444.3 (25)			
$(\Delta)^b L_{JK} \times 10^{12}$	-52.75 (51)			
$(\Delta)^b L_{JJK} \times 10^{12}$	-7.901 (32)			
$(\Delta)^b L_J \times 10^{15}$	-3.10 (17)	-2.64 (20) ^d	-5.28 (40) ^d	-5.28 (40) ^d
$(\Delta)^b P_{JK} \times 10^{18}$	552. (68)			
$(\Delta)^b P_{JJK} \times 10^{18}$	55.3 (22)			
$(\Delta)^b eQq$	-4.223 08 (107)	-0.038 7 (19) ^d	-0.077 4 (38) ^d	-0.077 4 (38) ^d
$eQq\eta$		0.1519 (113)		
$(\Delta)^b C_{bb} \times 10^3$	1.845 (90)			
$(\Delta)^b (C_{aa} - C_{bb}) \times 10^3$	-1.15 (31)			
$A\zeta$		138 656.195 (73)		138 656.042 (102)
η_K		10.332 9 (72)		10.405 1 (98)
η_J		0.390 469 (7)		0.394 512 (6)
$\eta_{KK} \times 10^6$		-834. (41) ^d		-834. (41) ^d
$\eta_{JK} \times 10^6$		-34.057 (64) ^d		-34.057 (64) ^d
$\eta_{JJ} \times 10^6$		-2.359 5 (24) ^d		-2.359 5 (24) ^d
$\eta_{JKK} \times 10^9$		2.59 (17) ^d		2.59 (17) ^d
$\eta_{JJK} \times 10^9$		0.508 5 (63) ^d		0.508 5 (63) ^d
q		17.798 438 (23)	^e	17.729 857 (138) ^e
$q_K \times 10^3$		-2.664 5 (111)	^e	-2.615 3 (89) ^e
$q_J \times 10^6$		-63.842 (14)	^e	-68.668 (31) ^e
$q_{JK} \times 10^9$		93.19 (53) ^d	^e	93.19 (53) ^{d,e}
$q_{JJ} \times 10^{12}$		311.5 (15)	^e	191.9 (26) ^e

^a Numbers in parentheses are one standard deviation in units of the least significant figures. Empty entries indicate parameter not applicable or not used in the fit. See section 4 for sign and value considerations.

^b Parameter X given for $v = 0$; $\Delta X = X_i - X_0$, with i representing an excited vibrational state.

^c All parameters given in units of megahertz, except for E_{vib} , which is given in units of inverse centimeters.

^d Ratio kept fixed in the fit for respective X or ΔX .

^e The parameter q and its distortion corrections connect levels with $\Delta K = \Delta l = 2$, see subsection 3.1 and section 4.

There is a $\Delta K = \Delta l = \pm 1$ Coriolis interaction between $v_4 = 1$ and $v_7 = 1$ which will be resonant at $K = 25$ or slightly higher for $v_7 = 1^{-1}$, too high to be observed easily. Another interaction occurs between $K = 6/7$ of $v_4 = 1$ and $K = 4/5$ of $v_7 = 1^{+1}$.

An anharmonic resonance occurs at $K = 7/8$ of $v_4 = 1$ and $v_8 = 3^{+3}$. Another, stronger anharmonic resonance occurs at $K = 7/8$ of $v_7 = 1^{+1}$ and $v_8 = 3^{+1}$. Inspection of preliminary energy levels involving these states suggest a plethora of further interactions among these states. Such interactions are also possible with states of the next polyad consisting of $v_4 = v_8 = 1$, $v_3 = 1$, $v_7 = v_8 = 1$, $v_6 = 1$, and $v_8 = 4$ [69].

4. Determination of spectroscopic parameters

Pickett's SPCAT and SPFIT programs [71] were used for prediction of the CH_3CN spectra and for fitting of the measured data. The programs were intended to be rather general, thus being able to fit asymmetric top rotors with spin and vibration-rotation interaction. They have evolved considerably with time because many aspects were not available initially [72, 73], in particular special considerations for symmetric or linear molecules or for higher symmetry cases. Spectra of symmetric top rotors with excitation of one [51, 74] or two quanta of a degenerate vibrational mode [51] were fit successfully with SPFIT. The l -doubling is implemented as $l \equiv 0, \pm 1 \pmod 3$, thus $l = \pm 2$ are identified with $l = \mp 1$. Hence, G_a , the $\Delta K = 0$ Coriolis interaction parameter between the l -components of $v_8 = 1^1$ and $v_8 = 2^2$ equals $-2A\zeta$ and $+4A\zeta$, respectively. In fact, the Coriolis term enters the Hamiltonian as $-2(A\zeta)_v kl$, see, e.g., Eq. (1) in Refs. [33] and [42]. The J - and K -dependent distortion parameters to G_a are usually defined without the factor of -2 , hence G_a^J equals η_J and $-2\eta_J$ for $v_8 = 1$ and $v_8 = 2$, respectively, and so on. Similarly, e.g., G_a and F_{ab} , are the coefficients of iJ_a and $(J_a J_b + J_b J_a)/2$, respectively. When they act as Coriolis parameters between two interacting states, the values from SPFIT often equal twice the values from other fitting programs. This is likely a consequence of the definition of G_i which is approximately equal to $2B_i \zeta_i$. F designates a Fermi (more generally anharmonic) interaction parameter. J or K as sub- or superscripts represent, as usual, corresponding distortion corrections. F_2 , G_{2i} , and F_{2ij} are the coefficients of operators which connect levels with ΔK higher by 2 than F , G_i , and F_{ij} , respectively, and with Δl modified accordingly to meet the symmetry condition $\Delta K - \Delta l \equiv 0 \pmod 3$. I.e. X_2 is defined for a parameter X as coefficient of $\{\mathbf{X}, \mathbf{J}_+^2 \mathbf{L}_+^2 + \mathbf{J}_-^2 \mathbf{L}_-^2\}$, with $\{a, b\} = (ab + ba)/2$ being the anticommutator. Interaction parameters between different vibrations are indicated by numbers in parentheses separated by a comma; the degree of excitation and the l quantum numbers are also given if necessary. The most important interaction parameter with $\Delta K > 0$ within a degenerate vibrational state of a strongly prolate symmetric rotor is q , also known as q_{22} because it connects levels differing in $\Delta K = \Delta l = 2$. It can have large effects in particular in low-lying bending modes, in which the magnitude of ζ often is close to the limiting case of 1.0, which will cause ample near-degeneracies. Scaling factors for q depending on the l -components and the vibrational excitation are not implemented in SPCAT and SPFIT.

The values are $q/2$ and $q/\sqrt{2}$ for $v_8 = 1$ and 2, respectively. We note that q and its distortion corrections are sometimes defined negatively. This applies to the studies on methyl cyanide [33, 42, 45, 69] and propyne [50, 51, 52]. Therefore, we inverted the signs and added footnotes to the appropriate tables. We also note an additional factor of 2 in the definitions of q_K , q_{JK} , etc. in the articles on propyne, causing the determined numerical values to be smaller than in SPFIT. However, none of these parameters were used in the IR studies of methyl cyanide relevant to our investigation [33, 42, 45, 69].

We determined rotational, centrifugal, and hyperfine structure (HFS) parameters of the ground state as common for all vibrational states. As most of the data, in particular the IR data, were not affected by HFS, all states were defined twice, with and without HFS. Vibrational changes $\Delta X = X_i - X_0$, where X represents a parameter and X_i the parameter in an excited vibrational state. This is very similar to several previous studies on CH_3CN [33, 42, 45, 69] and rather convenient because vibrational corrections ΔX are usually small with respect to X especially for lower order parameters X . In addition, this offers the opportunity to constrain vibrational corrections to $v_8 = 2$ to twice those of $v_8 = 1$, thus reducing the amount of independent spectroscopic parameters further.

The degenerate vibrational states required $A\zeta$ and q to be used in the fit with centrifugal distortion corrections as far as needed. Several of the distortion corrections between $v_8 = 1$ and 2 were constrained to be identical. Fitting of our new low frequency data of $v_8 = 1$ required $eQq\eta$ to be used in the fit, the quadrupole parameter equivalent to q . We assumed that the value determined by SPFIT is $+eQq\eta/4$.

Starting values for the parameters of $v = 0$, $v_8 = 1$, and $v_8 = 2$, were taken from Refs. [27, 33, 42]. Parameters describing the triad around 1000 cm^{-1} , $v_4 = 1$, $v_7 = 1$, and $v_8 = 3$, were taken from Ref. [45] and were kept fixed with the exception of one rotational interaction parameter between $v_8 = 2$ and $v_4 = 1$.

Parameters were introduced newly to the fit, or a constraint was released, if the rms error (as the measure of the quality of the fit) was reduced by an amount deemed significant and if the new parameter was useful in the Hamiltonian and if the resulting value(s) appeared to be reasonable. The latter aspect was a difficult one because vibrational corrections to some higher order parameters were of similar magnitudes as the parameters themselves and because some vibrational changes to $v_8 = 2^0$ or $v_8 = 2^2$ were rather different from twice the corresponding vibrational change to $v_8 = 1$. In most instances, care was taken to search for the largest reduction of the rms error. In almost all instances, this led to parameters determined with high significance, i.e. their values were more than five times larger in magnitude than their uncertainties. In few cases, parameters not determined with high significance were kept in the fit because the uncertainties were of similar magnitudes as those for other vibrational states.

Concerning the interactions between $v_8 = 1$ and 2, we determined J distortion corrections F_J and F_{JJ} to the main Fermi term F , whereas F_K was not determined with significance. We estimated its value by assuming $F_K/F_J \approx A/B$. Once both F_2 parameters were well-determined, their ratio was essentially 2

Table 4: Spectroscopic parameters or differences Δ thereof ^{a,b} (cm⁻¹, MHz)^c of methyl cyanide within vibrational states $v_4 = 1$, $v_7 = 1$, and $v_8 = 3$ taken from Ref. [45].

Parameter X	$v_4 = 1$	$v_7 = 1$	$v_8 = 3^1$	$v_8 = 3^3$
E_{vib}^c	920.290 03	1 041.854 71	1 077.786 3	1 122.15
$\Delta(A - B)$	-166.205	889.440	-302.030	-347.76
ΔB	-46.148 22	-5.734 13	80.294 85	81.289
$\Delta D_K \times 10^3$	-31.63	149.333	-22.8 ^d	-22.8 ^d
$\Delta D_{JK} \times 10^3$	7.165	0.973 1	3.504 6 ^d	3.504 6 ^d
$\Delta D_J \times 10^6$	-5.702 05	14.001	462. ^d	462. ^d
$\Delta H_K \times 10^6$	-22.24			
$\Delta H_{KJ} \times 10^6$	-13.895			
$\Delta H_{JK} \times 10^9$	315.68	-148.1		
$A\zeta$		66 663.668	138 665.87	138 527.8
η_K		7.224 8	11.013 ^d	11.013 ^d
η_J		0.078 153	0.401 54 ^d	0.401 54 ^d
$\eta_{KK} \times 10^6$		64.8		
$\eta_{JK} \times 10^6$		-50.34		
$\eta_{JJ} \times 10^6$		2.385	7.285 ^d	7.285 ^d
q		4.763 4	17.683 ^{d,e}	17.683 ^{d,e}
$q_J \times 10^6$		-10.85	-75.79 ^{d,e}	-75.79 ^{d,e}

^a Empty entries indicate parameter not applicable or not used in the fit. Signs of q and q_J altered, see section 4.

^b Parameter difference $\Delta(X) = X_i - X_0$, given for rotational and distortion parameters (i represents an excited vibrational state).

^c All parameters given in units of megahertz, except for E_{vib} , which is given in units of inverse centimeters.

^d Ratio kept fixed in the fit for parameter X or ΔX .

^e The parameter q and its distortion corrections connect levels with $\Delta K = \Delta l = 2$, see subsection 3.1 and section 4.

within the uncertainties. Therefore, the ratio was constrained to reduce the number of independent parameters. No distortion corrections could be determined for the Fermi terms associated with the interactions between $v_8 = 2$ and 3, so the J and K distortion corrections were fixed to appropriate multiples of the values determined for the $v_8 = 1/2$ resonance. Trial fits with F_2 yielded values or constraints which were in some cases much smaller than those in the $v_8 = 1/2$ resonances with essentially no effect on the quality on the fit. These parameters were omitted from the final fits.

The parameter f_{44} (for $v_8 = 2$) had little effects on the rms error in most trial fits and was not determined with significance. It seemed to be determined in some trial fits, but its magnitude appeared to be too large with respect to that of $v_{10} = 2$ of propyne [51]. Moreover, changes in some of the remaining spectroscopic parameters were deemed too large in these cases. Therefore, f_{44} was omitted from the final fit.

5. Discussion

Our present ground state spectroscopic parameters are very close to those of our previous study [27]; differences outside three times the experimental uncertainties occur only in D_J , H_J , and L_J . The experimental data associated with the $v = 0/v_8 = 1$ perturbation help to constrain the K level structure in the ground vibrational state, additional contributions may come from the data involving the $v_8 = 1/v_8 = 2$ perturbations. As a consequence, we obtained a precisely determined value for H_K as (164.6 \pm 6.6) Hz. The agreement with (156 \pm 72) Hz from an earlier study [32] is good, though the large uncertainty in the earlier study limits the value of the agreement somewhat.

Table 5: Interaction parameters^a (MHz) between vibrational states of methyl cyanide.

Parameter	Value
$F_2(0, 8^1) \times 10^3$	-70.897 (27)
$F(8^{\pm 1}, 8^{2,\mp 2})$	53 157.7 (33)
$F_K(8^{\pm 1}, 8^{2,\mp 2})$	-6. ^{b,c}
$F_J(8^{\pm 1}, 8^{2,\mp 2}) \times 10^3$	-369.89 (44) ^b
$F_{JJ}(8^{\pm 1}, 8^{2,\mp 2}) \times 10^6$	1.681 (87) ^b
$F_2(8^{\pm 1}, 8^{2,0}) \times 10^3$	-65.491 (24) ^d
$F_2(8^{\pm 1}, 8^{2,\pm 2}) \times 10^3$	-130.982 (48) ^d
$F_{ac}(8^{2,\pm 2}, 4^1) [2W_{488}]$	8.736 2 (21)
$F_{2ac}(8^{2,0}, 4^1) \times 10^6$	-7.98 (16)
$F(8^{\pm 2}, 7^{\mp 1}) [W_{788}]$	45 170.8 ^e
$F(8^{2,\pm 2}, 8^{3,\mp 1})$	77 208. (93)
$F(8^{2,0}, 8^{3,3})$	91 509. (131)
$G_b(4, 7) [2W_{47}]$	909. ^e
$F_{bc}(4, 7) [2W_{47}]$	-1.84 ^e
$F(4, 8^{3,\pm 3}) [W_{4888}]$	11 430. ^e
$F(7^{\pm 1}, 8^{3,\pm 1}) [W_{7888}]$	50 129.2 ^e
$G_a(7^{\pm 1}, 8^{3,\pm 1}) [2W_{7888}^k]$	-2 239.1 ^e

^a Alternative designations from Ref. [45] given in brackets. Numbers in parentheses after the interaction parameter designate the vibrational states separated by a comma, see also section 4. Numbers in parentheses after the values are one standard deviation in units of the least significant figures.

^b J and K distortion corrections to $F(8^{2,\pm 2}, 8^{3,\mp 1})$ and $F(8^{2,0}, 8^{3,3})$ kept fixed to $\sqrt{2}$ and $\sqrt{3}$, respectively, times the corresponding $F(8^{\pm 1}, 8^{2,\mp 2})$ value; see, e.g., Ref. [51].

^c Estimated assuming $F_K/F_J \approx A/B$, see section 4.

^d Ratio constrained, see section 4.

^e Kept fixed to values from Ref. [45].

Table 6: Spectroscopic parameters or differences Δ thereof ^{a,b} (cm⁻¹, MHz)^c of methyl cyanide involving $v_8 = 1$ and 2 in comparison to previous values.

Parameter X	$v_8 = 1$		$v_8 = 2^0$		$v_8 = 2^2$	
	Present	Ref. [33]	Present	Ref. [42]	Present	Ref. [42]
E_{vib}^c	365.024365 (9)	365.015965 (12)	716.75042 (13)	716.70441 (1)	739.14823 (6)	739.12907 (3)
ΔA	-88.400 (26)	-90.079 (36)	-133.347 (18)	-140.069 (39)	-205.453 (122)	-210.826 (96)
ΔB	27.53028 (5)	27.53255 (33)	54.05732 (11)	54.07405 (54)	54.50273 (7)	54.50740 (90)
$\Delta D_K \times 10^3$	-11.46 (48)	-18.11 (81)	-20.2 (3)	-12.3 (10)	-7.5 (16)	20.1 (42)
$\Delta D_{JK} \times 10^3$	0.9875 (6)	0.8739 (129)	1.6755 (25)	1.3316 (69)	1.809 (2)	1.790 (39)
$\Delta D_J \times 10^6$	95.599 (17)	93.436 (69)	216.32 (4)	218.41 (18)	189.16 (3)	182.93 (24)
$\Delta H_K \times 10^6$	14.9 (22)	-215.0 (51)		-384.3 (72)		
$\Delta H_{KJ} \times 10^6$	0.034 (2)	0.908 (114)	0.150 (24)		0.025 (6)	
$\Delta H_{JK} \times 10^9$	2.59 (6)	35.74 (171)	17.71 (34)		-0.37 (21)	
$\Delta H_J \times 10^{12}$	315.3 (30)		200.7 (63)		627.8 (59)	
$\Delta L_J \times 10^{15}$	-2.64 (20) ^d		-5.28 (40) ^d		-5.28 (40) ^d	
$A\zeta \times 10^{-3}$	138.65620 (7)	138.64715 (7)			138.65604 (10)	138.63733 (16)
η_K	10.333 (7)	10.448 (93)			10.405 (10)	10.448 (9)
η_J	0.390469 (7)	0.389679 (66)			0.39451 (1)	0.39420 (18)
$\eta_{KK} \times 10^6$	-834. (41) ^d	-453. (174)			-834. (41) ^d	
$\eta_{JK} \times 10^6$	-34.06 (6) ^d	-40.77 (90)			-34.06 (6) ^d	-46.8 (30)
$\eta_{JJ} \times 10^6$	-2.3595 (24) ^d	-2.6750 (111)			-2.3595 (24) ^d	-2.845 (45)
$\eta_{KKK} \times 10^6$		15.23 (90)				
$\eta_{JKK} \times 10^9$	2.59 (17) ^d				2.59 (17) ^d	
$\eta_{JJK} \times 10^9$	0.509 (6) ^d				0.509 (6) ^d	
q	17.79844 (2)	17.79776 (11)	<i>e</i>	<i>e</i>	17.72986 (14) ^e	17.70068 (51) ^e
$q_K \times 10^3$	-2.6645 (111)		<i>e</i>	<i>e</i>	-2.6153 (89) ^e	
$q_J \times 10^6$	-63.84 (1)	-63.80 (16)	<i>e</i>	<i>e</i>	-68.67 (3) ^e	-66.84 (17) ^e
$q_{JK} \times 10^9$	93.19 (53) ^d		<i>e</i>	<i>e</i>	93.19 (53) ^{d,e}	
$q_{JJ} \times 10^{12}$	312. (2)	324. (57)	<i>e</i>	<i>e</i>	192. (3) ^e	

^a Numbers in parentheses are one standard deviation in units of the least significant figures. Empty entries indicate parameter not applicable or not used in the fit. See section 4 for sign and value considerations.

^b Parameter difference $\Delta(X) = X_i - X_0$ given for rotational and distortion parameters (i represents an excited vibrational state). Signs were adjusted to present definition for data from Refs. [33, 42]. In all instances, previous values refer to the combined fits of IR and rotational data.

^c All parameters given in units of megahertz, except for E_{vib} , which is given in units of inverse centimeters.

^d Ratio kept fixed in the fit for respective X or ΔX .

^e The parameter q and its distortion corrections connect levels with $\Delta K = \Delta l = 2$, see subsection 3.1 and section 4.

Our previous estimate of 51 Hz, derived by scaling D_K with the D_K/A ratio as an order-of-magnitude estimate, was apparently too small.

As can be seen in **Table 6**, the extensive new rotational data pertaining to $v_8 = 1$ and 2 allowed us to determine spectroscopic parameters or changes thereof up to higher order than previously [33, 42]. The uncertainties of many lower order parameters were improved in addition.

The previous and present spectroscopic parameters agree reasonably at lower order and particularly for those with dependence on J . The agreement deteriorates for higher order parameters, especially those dependent on K . Consideration of the resonances between $v_8 = 1$ and 2 as well as between $v_8 = 2$ and 3 affected the purely K -dependent parameters significantly, very similar to the case of propyne, where the Fermi interactions involving the equivalent ν_{10} band occur at slightly lower K [50, 51]. As in the case of propyne, we did not require η_{KKK} to be included in the fit. Comparison of the various η parameters from our study as well as those from Ref. [33] suggests that the magnitude of η_{KKK} reported in that study is at least a factor of 100 too large. Similarly, the previous ΔH_K value of $v_8 = 1$ in CH_3CN was -215 ± 5 Hz [33], whereas our value is 14.9 ± 2.2 Hz; in addition, a value of -384 ± 7 Hz was reported for $v_8 = 2^0$, in Ref. [33], whereas we obtained a value similar to that in $v_8 = 1$ in trial fits. However, the uncertainties were only slightly smaller than the values, so $\Delta H_K(v_8 = 2^0)$ was omitted in the final fit.

The vibrational changes in distortion parameters from $v = 0$ to $v_8 = 1$ in **Table 3** are often small, though not in all instances. The rather large changes in H_J and L_J are common in these types of molecules. In fact, the changes of H_J upon excitation of each of the two low-lying bending modes are even larger in magnitude in propyne [51]. Excitation of the CCN bending mode leads also to large changes in H of HC_3N [77], DC_3N [78], or other, similar molecules.

The changes of rotational and centrifugal distortion parameters from ground state values in $v_8 = 2^0$ and $v_8 = 2^2$ are roughly twice those in $v_8 = 1$, though pronounced deviations occur already in ΔA (or $\Delta(A - B)$). The similarity to related values in propyne [50, 51] and methyl isocyanide [75, 76] may indicate that the low energies of the vibrations lead to substantial changes in ΔA with different l quantum number. Changes upon excitation in v_8 are closer to the ratio of 2 in ΔB , but deviations increase, in particular for the changes in the ΔH values. These deviations may be caused by the incomplete analysis of the $v_8 = 2/3$ or other resonances.

The $A\zeta$ values are essentially identical between $v_8 = 1$ and $v_8 = 2^2$. The changes in η_K and η_J are rather small, and higher order parameters were constrained to the same values. The q values are rather similar also, whereas the q_{JJ} values differ considerably, possibly again a consequence of a not or not fully accounted interaction.

With respect to the perturbation of the ground vibrational state by $v_8 = 1$, Šimečková et al. [41] calculated perturbations of about 45 MHz and 36 MHz for the two ground state transitions involving $J = 43$ and $K = 14$, this is about 20 MHz larger than actually observed. The difference can be traced al-

most entirely to the energy differences between the two perturbing levels at $J = 43$, which was estimated to be 0.013 cm^{-1} [41] whereas our calculations yield 0.025 cm^{-1} . Our value of $F_{ac}(8^{2,\pm 2}, 4^1)/2$, 4.368 ± 0.001 MHz agrees quite well with $|w_{488}| = 4.101 \pm 0.042$ MHz from the combined analysis of IR and rotational data or with 4.170 ± 0.042 MHz from an analysis of IR data only [45]. We should point out that the sign cannot be determined from the fit of the line frequencies, but it may be determinable from relative intensities in the rotational or IR spectra. The value given in Ref [45] was actually negative; we will try to determine the sign in our subsequent study which will include $v_4 = 1$ data. Other interaction parameters between different vibrational states were either determined for the first time here or were not released in our fits.

Selected low order spectroscopic parameters of CH_3CN have been compared in **Table 7** with corresponding values of CH_3CCH and CH_3NC . The data of CH_3CN and of CH_3CCH are comparatively similar, while those of CH_3NC with its much lower bending mode are substantially different. The Fermi parameters between the first and the second bending state reflect these changes, and the values appear to scale approximately with q_b .

$F(8^{2,\pm 2}, 8^{3,\mp 1})$ and $F(8^{2,0}, 8^{3,3})$ should ideally differ from $F(8^{\pm 1}, 8^{2,\mp 2})$ by factors of $\sqrt{2}$ and $\sqrt{3}$, respectively, see, e.g., Ref. [51]. The experimental values of 77208 ± 93 MHz and 91509 ± 131 MHz compare quite well with the expected values of 75192 MHz and 92091 MHz, respectively, in particular considering that there are no $v_8 = 3$ transition frequencies in the fit, and that we used spectroscopic parameters from a previous study for this state [45].

The lack of assignments in $2\nu_8 - \nu_8$ may limit the accuracies of the predicted line positions in the $2\nu_8^{+2} - \nu_8^{+1}$ sub-band at higher K . However, our data on the perturbations between $v_8 = 1$ and 2 may in part compensate these limitations.

6. Expanded database for CH_3CN line parameters

Compilations of CH_3CN molecular line parameters are greatly expanded by including new or updated predictions of spectra within or between the three states of the present investigation. Predictions of rotational and infrared spectra for astrophysical applications will be available in the catalog section⁴ of the CDMS [1, 2] and that of JPL⁵ [79]. The parameter, line, fit and other auxiliary files are available online as supplements with the publication and in the Cologne Spectroscopy Data section⁶ of the CDMS. For remote sensing of, e.g., Titan, the new compilation will also be provided to HITRAN [64] and GEISA [80] with N_2 -pressure and self-broadening information included. For this, the extensive pressure broadening CH_3CN measurements of P and R branch transitions of ν_4 [68] were modeled by extending *empirical expressions* applied to ammonia [81, 82]. The resulting polynomial has no physical meaning,

⁴<https://cdms.astro.uni-koeln.de/classic/entries/>

⁵<http://spec.jpl.nasa.gov/ftp/pub/catalog/catdir.html>

⁶<https://cdms.astro.uni-koeln.de/classic/predictions/daten/CH3CN/CH3CN/>

Table 7: Comparison of selected low order spectroscopic parameters or differences Δ thereof^a (cm⁻¹, MHz)^b of methyl cyanide, methyl acetylene, and methyl isocyanide in their ground and lowest bending states b.

Parameter X	CH ₃ CN ^c	CH ₃ CCH ^d	CH ₃ NC ^e
A_0	158 099.	159 139.	157 308.
B_0	9 198.9	8 545.9	10 052.8
$D_{K,0} \times 10^3$	2 831.	2 939.	2 567.
$D_{JK,0} \times 10^3$	177.4	163.4	227.5
$D_{J,0} \times 10^6$	3 808.	2 908.	4 692.
E_b^b	365.0	330.9	267.3
$\Delta^a A_b$	-88.3	-61.9	0.6
$\Delta^a B_b$	27.53	23.94	39.07
$A\zeta_b$	138 656.	141 919.	146 007.
$\eta_{K,b}$	10.34	10.94	10.13
$\eta_{J,b} \times 10^3$	390.5	334.8	587.6
q_b	17.80	16.79	6.94
$F(b^{\pm 1}, b^{2, \mp 2})$	53 169.	51 745.	23 808.

^a Parameter X given for $v = 0$; $\Delta(X) = X_i - X_0$, with i representing an excited vibrational state.

^b All parameters given in units of megahertz, except for E_b , which is given in units of inverse centimeters.

^c This work.

^d Ref. [51]; sign of q adjusted, see section 4.

^e Refs. [53, 75, 76]; sign of q adjusted, see section 4.

but it does provide reasonable estimates of Lorentz broadening coefficients for CH₃CN transitions.

6.1. Pressure-broadened half widths

The measured N₂- and self-broadened halfwidths (HWHM) for CH₃CN in the ν_4 band at 920 cm⁻¹ vary smoothly with respect to their quantum numbers J and K [68]. Patterns become more obvious when the widths are examined as a function of values related to the transition quantum numbers; these are defined as $J_m = J'' + 1$, $K_m = K'' + 1$ for R branch transitions, and $J_m = J''$, $K_m = K''$ for P and Q branch transitions. Examples are shown in **Table 8**. N₂- and self-broadened half widths were derived for our CH₃CN compilation by modeling the reported widths for the ν_4 band [68] according to the expression in Eq. 1; the polynomial coefficients are listed in **Table 9**.

$$\gamma(J_m, K_m) = \sum_{i,j} a_{ij} J_m^i K_m^j \quad (1)$$

The measured widths were weighted in the fit by the inverse of the square of the reported uncertainties [68]. The impact of a few outliers (at very low and higher J_m and K_m) were suppressed as an important side-effect. As shown in **Fig. 10**, the calculated values using Eq. 1 (solid curves) reproduce the pattern of the measured data points within 5% for most of the transitions. The full lists are given in Supplemental files 5_Width(N2)_CH3CN_v4.out and 6_Width(Self)_CH3CN_v4.out. Finally, we adopted this expression to estimate N₂- and self-broadened half widths for the three IR bands, ν_8 , $2\nu_8 - \nu_8$, $2\nu_8$ of the present study. Because the experimental ν_4 lines, for which line widths were determined, have $J_m \lesssim 50$, we limited the line-widths at 296 K to:

Table 8: Examples of measured^a (obs.; with uncertainties unc. in percent in parentheses) and modeled (calc.; with percent difference in parentheses) ν_4 self-broadened line widths (cm⁻¹/atm) of methyl cyanide at 296 K with frequencies (cm⁻¹) and quantum numbers.

J', K'	J'', K''	Frequency	Self-broadened width		J_m, K_m
			obs. (%unc.)	calc. (%diff.)	
2, 1	3, 1	918.43291	1.358 (0.7)	1.2588 (7.88)	3, 1
3, 1	2, 1	922.10564	1.307 (1.0)	1.2588 (3.83)	3, 1
7, 6	8, 6	915.09499	1.071 (0.7)	1.0798 (-0.82)	8, 6
8, 6	7, 6	924.88187	1.039 (0.9)	1.0798 (-3.78)	8, 6
18, 9	19, 9	907.67687	1.722 (0.5)	1.6957 (1.55)	19, 9
19, 9	18, 9	930.89430	1.711 (0.7)	1.6957 (0.90)	19, 9
22, 7	23, 7	905.12404	1.773 (0.4)	1.7497 (1.33)	23, 7
23, 7	22, 7	933.24349	1.821 (0.6)	1.7497 (4.07)	23, 7
28, 0	29, 0	901.25148	1.583 (0.5)	1.5807 (0.15)	29, 0
29, 0	28, 0	936.73087	1.528 (0.5)	1.5807 (-3.33)	29, 0
40, 4	41, 4	892.54529	0.896 (0.7)	0.8737 (2.55)	41, 4
41, 4	40, 4	942.65553	0.823 (0.8)	0.8737 (-5.81)	41, 4

^a The measured values are taken from Ref. [68].

Table 9: Coefficients to model the N₂- and self-broadened widths of the ν_4 band of CH₃CN using Eq. 1.

Coefficients	N ₂	Self
a_{00}	0.171	1.621
$a_{10} \times 10^3$	4.731	-172.4
$a_{01} \times 10^3$	-9.962	-82.99
$a_{20} \times 10^3$	-0.7152	26.32
$a_{11} \times 10^3$	1.704	18.89
$a_{02} \times 10^3$	-1.233	-14.42
$a_{30} \times 10^3$	0.03155	-1.28
$a_{21} \times 10^3$	-0.1068	-1.296
$a_{12} \times 10^3$	0.1361	0.7237
$a_{03} \times 10^3$	0.002544	0.1243
$a_{40} \times 10^6$	-0.6094	24.83
$a_{31} \times 10^6$	2.643	-34.66
$a_{22} \times 10^6$	-3.849	-10.5
$a_{13} \times 10^6$	-1.819	9.179
$a_{50} \times 10^9$	4.388	-171.6
$a_{41} \times 10^9$	-22.42	-319.8
$a_{32} \times 10^9$	33.41	66.99
$a_{23} \times 10^9$	37.85	-430.9

$0.12 < \gamma_{N_2} < 0.2 \text{ cm}^{-1}/\text{atm}$ for N_2 -broadened half-widths, and

$0.2 < \gamma_{\text{self}} < 2.0 \text{ cm}^{-1}/\text{atm}$ for self-broadened half-widths.

We adopted one single value, 0.72, from the HITRAN2012 database [64] for the temperature dependence exponent of the N_2 -broadened Lorentz widths. This value is based on the measurements of the $J = 12 - 11$, $K = 3$ transition in the ground vibrational state [83].

6.2. N_2 -pressure-induced frequency shifts

Rinsland et al. [68] also reported N_2 -pressure-induced frequency shifts for CH_3CN in its ν_4 band, but no distinctive pattern is seen for pressure shifts with respect to J_m and K_m . However, there was a significant correlation between measured line widths and frequency shifts, as presented in their Fig. 11 of Ref [68]. We use this observation in our present work to obtain frequency shift estimates for the transition of ν_8 , $2\nu_8 - \nu_8$, and $2\nu_8$. First, we apply a scaled set of the *measured* N_2 -induced frequency shifts for ν_4 transitions for all other CH_3CN transitions having the same set of quantum numbers, J'' , K'' , J' , and K' . However, not all transitions predicted in the present work have counter parts among the measured ν_4 transitions; we assume for these cases that their N_2 -induced frequency shifts $\delta_4(J_m, K_m)$ can be estimated from their pressure-broadened line widths using the *width-shift relationship* proposed by Rinsland et al. [68]. Based on Fig. 11 of Ref. [68], a simple expression was derived as

$$\delta_4(J_m, K_m) = [\gamma(J_m, K_m)_{N_2} - 0.15]/15 \quad (2)$$

One commonly assumes that the overall magnitudes of the pressure shifts increase with increasing vibrational band centers [84]. We estimated the pressure shifts $\delta_i(J_m, K_m)$ of ν_i from those of the ν_4 band according to

$$\delta_i(J_m, K_m) = \delta_4(J_m, K_m) \times (E_{\text{vib}}(i)/E_{\text{vib}}(4)) \quad (3)$$

with $E_{\text{vib}}(i)$ being the vibrational energy of ν_8 , $2\nu_8 - \nu_8$, or $2\nu_8$ and $E_{\text{vib}}(4)$ that of ν_4 .

The resulting N_2 -pressure shifts for the ν_8 , $2\nu_8 - \nu_8$, and $2\nu_8$ bands are presented in **Fig. 11a**. The width-shift relationships resemble those observed in the ν_4 band [68]. The distribution with respect to K_m , shown in **Fig. 11b**, does not show a distinctive pattern, as was the case for ν_4 [68].

6.3. Intensities of ν_8 , $2\nu_8 - \nu_8$, and $2\nu_8$

We evaluated transition dipole moments for each of the three IR bands of the present study. There are some indications that additional correction terms may lead to a better intensity modelling, especially in the cases of ν_8 and $2\nu_8 - \nu_8$. More sophisticated intensity modelling is beyond the aim of the present study, in particular because the quality of the IR spectrum in the ν_8 region is not sufficient for such modelling.

6.3.1. ν_8

Our simulations on a line-by-line basis indicate that the ν_8 spectrum is represented reasonably well by a transition dipole moment of $\mu(8_0^1) = 0.043 \text{ D}$.

We calculated partition function values by summing over all states up to $J = 99$, $K = 25$, and $\nu_8 = 3$. While the rotational contribution to the partition function at 296 K is well converged, the vibrational contribution is slightly truncated. We have estimated the truncation error classically, assuming a rigid rotor as well as a harmonic oscillator. This approximation in itself introduces some error, but we estimate it to be small with respect to the truncation error. The truncation leads to a corrections of 1.408 %. The resulting intensity, integrated over all lines, is given in Table 10. The integrated intensity in the entire ν_8 region, i.e. with all hot bands, was estimated by multiplying the calculated ν_8 intensity with the vibrational partition factor of 1.5006. An additional increase of the integrated intensity by 2.677 % was required because of the contributions of other isotopic species. There is a convention in HITRAN that intensities refer to a sample in natural isotopic composition. The resulting intensity is also given in Table 10. Our value is about half way between two lower values [67, 85] and one higher one [86].

6.3.2. $2\nu_8 - \nu_8$

We assumed transition dipole moments of $\mu(8_l^{2'}) = 0.043 \text{ D}$ for both $l = 0$ and $l = \pm 2$ components as for the cold band. This approximation appears to be valid well within the estimated intensity uncertainties of 4 %, see subsection 6.3.3. **Fig. 12** shows the calculated contributions of the ν_8 and $2\nu_8 - \nu_8$ sub-branches.

As can be seen in Table 10, slightly more than 10 % of the total intensity in the ν_8 region of methyl cyanide are caused by IR transitions originating in states higher than $\nu_8 = 1$. Finding unblended lines among these will be difficult because this was already not easy for $2\nu_8^2 - \nu_8$.

6.3.3. $2\nu_8$

The relative precisions of measured intensities were checked using spectral comparisons. For example, the 731–732 cm^{-1} interval in **Fig. 2** has 181 retrieved features with intensities ranging from 1.3×10^{-6} to $7.3 \times 10^{-3} \text{ cm}^{-2}/\text{atm}$ at 293 K. Since the sample was a mixture, line intensities of the two calibration species were measured to obtain the partial pressures of OCS (3 %) and CO_2 (0.4 %). In addition, selected strong CH_3CN lines of the ν_4 band were retrieved and compared to values reported in Table 2 of Ref. [68]; this indicated that the retrieved intensities in the $2\nu_8$ region should be increased by 4 %; with these calibrations, the overall (relative) precision is thought to be 0.0002 cm^{-1} for line centers and 5 % for intensities. However, worse precisions are expected for blended features separated by less than 0.0025 cm^{-1} and for ~ 5100 absorptions weaker than $9 \times 10^{-5} \text{ cm}^{-2}/\text{atm}$ at 293 K.

Our simulations indicate that the spectrum is represented well by a transition dipole moment of $\mu(8_0^{20}) = 0.030 \text{ D}$. See subsection 2.3 for considerations concerning the experimental intensities and subsection 6.3.1 for a remark on the partition

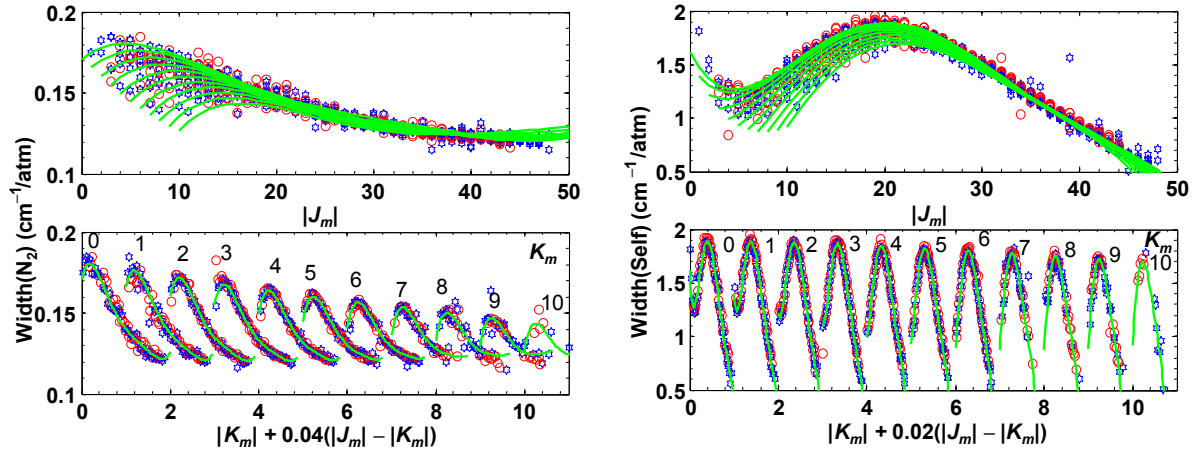


Figure 10: Model fit of measured N_2 - (left) and self-broadening (right) Lorentz width coefficients in $\text{cm}^{-1}/\text{atm}$ at 296 K for the ν_4 band [68]. Many of the widths are well reproduced by Eq. 1. The solid lines in green represent calculated values while observed widths of P and R branch transitions are indicated by circles in red and hexagons in blue, respectively. In the lower parts, the J_m and K_m values are combined for the horizontal axes to reveal the variation patterns.

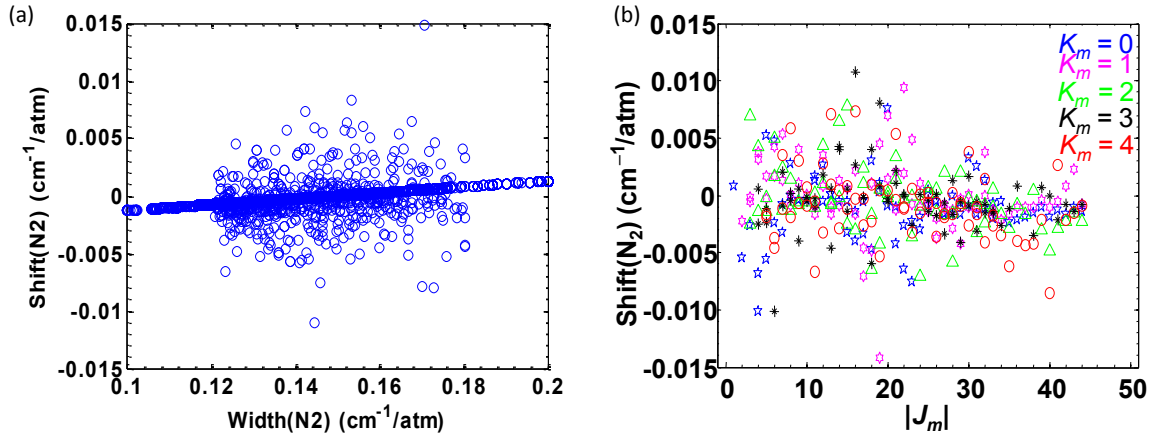


Figure 11: (a) N_2 -pressure induced frequency shifts derived for the ν_8 region. The scattered values are adapted from the ν_4 measurements [68] while the well-aligned ones are based on Eqs. 2 and 3. (b) The derived N_2 -shifts of IR bands in the present study are presented for selected K_m (see text for details).

Table 10: Derived transition dipole moments μ (D) and integrated intensities (10^{-19} $\text{cm}/\text{molecule}$) of IR bands of $^{12}\text{CH}_3^{13}\text{C}^{14}\text{N}$ calculated from the present Hamiltonian model with respect to the transition dipole moments (Calc. (μ)) and apparent intensities in the l components (Calc. (l)) together with the extrapolated and experimentally measured total integrated intensities from previous studies.

IR band	μ	Integrated intensity					
		Calc. (μ)	Calc. (l)	Ref. [66]	Ref. [67]	Ref. [85]	Ref. [86]
$\nu_8 = 1 - 0$	0.043	1.483	1.483				
$\nu_8 = 2^0 - 1$	0.043	0.251	0.244				
$\nu_8 = 2^2 - 1$	0.043	0.257	0.264				
ν_8 region		2.285 ^a		1.77 (4)	1.81	2.79 (28)	
$\nu_8 = 2^0 - 0$	0.030	1.674	1.585				
$\nu_8 = 2^2 - 0$	0.000	0.000	0.090				
$2\nu_8$ region		2.580 ^a		2.63 (18)	2.50 (4)	3.43	

^a Estimated from the cold bands, $\nu_8 = 1 - 0$ and $2 - 0$, by multiplication with the vibrational partition factor 1.5006 at 296 K. An increase of 2.677% from other isotopic species was considered.

function and on the sample composition. The integrated intensity from our Hamiltonian model is 1.75×10^{-19} cm/molecule for the $2\nu_8$ band alone. Only a small portion of this intensity, $\sim 1.0 \times 10^{-21}$ and $\sim 8.1 \times 10^{-21}$ cm/molecule is transferred through the q_{22} interaction to the $\nu_8 = 2^{+2}$ and $\nu_8 = 2^{-2}$ components, respectively.

About 23 % of the $2\nu_8$ intensity at 296 K should arise from the $3\nu_8 - \nu_8$ and $\nu_7 - \nu_8$ bands. A preliminary analysis suggests that a fair fraction of these lines will not be blended by other lines of similar or higher intensity. Even some of the lines originating in states higher than $\nu_8 = 1$, which account for slightly more than 10 % of the band intensity, may be usable for spectroscopic analyses. More work is needed to clarify these aspects.

7. Conclusions

Our extensive and accurate rotational data in vibrational states $\nu_8 \leq 2$ resulted in a parameter set extended to higher order and to mostly improved uncertainties at lower order. More importantly, these data enabled the first treatments of resonances between $\nu_8 = 1$ and 2, as well as between $\nu_8 = 2$ and 3. These resonances cause considerable perturbations in the rotational spectrum, and their successful treatments affect the spectroscopic parameters, in particular the K -dependent ones and those at higher order. We have also performed the first in-depth analysis of the resonant interaction between the ground vibrational state and $\nu_8 = 1$.

The analysis of a new $2\nu_8$ spectrum was an additional important contribution to the determination of spectroscopic parameters. Only about half of the roughly 12000 lines in the $2\nu_8$ spectrum are caused by the cold band. A large fraction of the unaccounted lines will be caused by the $3\nu_8 - \nu_8$ hot band and the $\nu_7 - \nu_8$ difference band. Accounting for these will require in depth analysis of the $\nu_7/3\nu_8$ band system aided by rotational transitions. A still considerable fraction of lines are due at least in part to bands originating in states higher than $\nu_8 = 1$.

We have derived line broadening and shift parameters for the ν_8 , $2\nu_8 - \nu_8$, and $2\nu_8$ bands from previously determined ν_4 values and have used this information to determine transition moments and intensities for the three IR bands. Individually measured transition intensities may lead to an improved intensity model.

We have generated an expanded database for CH_3CN line parameters for transitions within each of the states $\nu_8 \leq 2$, mainly for radio astronomical observations, in particular with ALMA and other interferometers. In addition, information is provided for transitions between these three states for remote sensing of, e.g., the atmosphere of Titan.

The inclusion of data pertaining to $\nu_4 = 1$ is almost complete. We intend to improve the data involving its interactions with other vibrational states. We want to study also the proposed strong, although distant, thus non-resonant Fermi interaction of $\nu_4 = 1$ with $\nu_8 = 2$ [43]. In addition, we have already made assignments in the $\nu_4 + \nu_8 - \nu_8$ IR band as well as for rotational transitions in $\nu_4 = \nu_8 = 1$, though these are often tentative. The spectra are perturbed, but, as it appears, only locally, such that a reasonable deperturbation may be possible by considering only spectroscopic parameters of the perturbing states.

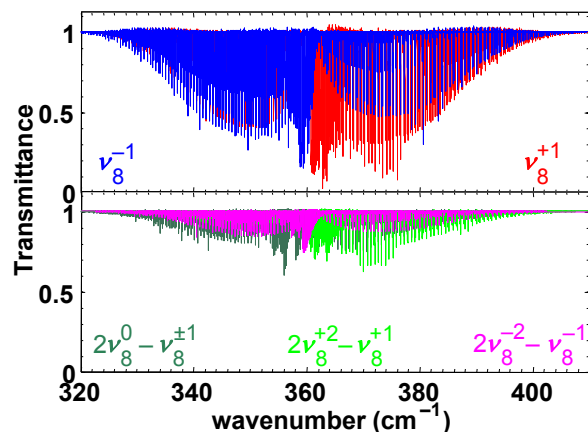


Figure 12: Contributions of each sub-band in the ν_8 region of CH_3CN showing their relative intensities for the cold band in the upper part and for the hot band in the lower part. Note: hot bands originating in states higher than $\nu_8 = 1$ contribute to $\sim 10\%$ to the total intensity.

Acknowledgements

H.S.P.M. is grateful to the Bundesministerium für Bildung und Forschung (BMBF) for financial support through project FKZ 50OF0901 (ICC HIFI *Herschel*) during part of the present investigation. The measurements in Köln were supported by the Deutsche Forschungsgemeinschaft (DFG) through the collaborative research grants SFB 494 initially, and later SFB 956, project area B3. The portion of this work, which was carried out at the Jet Propulsion Laboratory, California Institute of Technology, was performed under contract with the National Aeronautics and Space Administration. The infrared spectra analyzed in the present study were recorded at the W.R. Wiley Environmental Molecular Sciences Laboratory, a national scientific user facility sponsored by the Department of Energy's Office of Biological and Environmental Research located at the Pacific Northwest National Laboratory (PNNL). PNNL is operated for the United States Department of Energy by the Battelle Memorial Institute under Contract DE-AC05-76RLO1830.

Appendix A. Supplementary Material

Supplementary data for this article are available on ScienceDirect (www.sciencedirect.com) and as part of the Ohio State University Molecular Spectroscopy Archives (http://library.osu.edu/sites/msa/jmsa_hp.htm). The parameter, line, and fit files along with a readme file are provided as well as the full tables of the modeling of N_2 - and self-broadening line widths of the ν_4 band along with an explanation file. Supplementary data associated with this article can be found, in the online version, at <http://dx.doi.org/10.1016/j.jms.2015.02.009>.

References

- [1] H.S.P. Müller, S. Thorwirth, D.A. Roth, G. Winnewisser, *Astron. Astrophys.* 370 (2001) L49–L52.

- [2] H.S.P. Müller, F. Schlöder, J. Stutzki, G. Winnewisser, *J. Mol. Struct.* 742 (2005) 215–227.
- [3] A. Belloche, R.T. Garrod, H.S.P. Müller, K.M. Menten, C. Comito, P. Schilke, *Astron. Astrophys.* 499 (2009) 215–232.
- [4] A. Belloche, R.T. Garrod, H.S.P. Müller, K.M. Menten, *Science* 345 (2014) 1584–1587.
- [5] D.R. Johnson, F.J. Lovas, C.A. Gottlieb, E.W. Gottlieb, M.M. Litvak, M. Guélin, P. Thaddeus, *Astrophys. J.* 218 (1977) 370–376.
- [6] K. Demyk, H. Mäder, B. Tercero, J. Cernicharo, J. Demaison, L. Margulès, M. Wegner, S. Keipert, M. Sheng, *Astron. Astrophys.* 466 (2007) 255–259.
- [7] H.S.P. Müller, A. Belloche, K.M. Menten, C. Comito, P. Schilke, *J. Mol. Spectrosc.* 251 (2008) 319–325.
- [8] E. Gibb, A. Nummelin, W.M. Irvine, D.C.B. Whittet, P. Bergman, *Astrophys. J.* 545 (2000) 309–326.
- [9] D.M. Mehringer, J.C. Pearson, J. Keene, T.G. Philips, *Astrophys. J.* 608 (2004) 306–313.
- [10] A.M. Daly, C. Bermúdez, A. López, B. Tercero, J.C. Pearson, N. Marcelino, J.L. Alonso, J. Cernicharo, *Astrophys. J.* 768 (2013) 81.
- [11] A. Belloche, H.S.P. Müller, K.M. Menten, P. Schilke, C. Comito, *Astron. Astrophys.* 559 (2013) A47.
- [12] N.J. Livesey, J.W. Waters, R. Khosravi, G.P. Brasseur, G.S. Tyndall, W.G. Read, *Geophys. Res. Lett.* 28 (2001) 779–782.
- [13] A. Kleinböhl, G.C. Toon, B. Sen, J.-F.L. Blavier, D.K. Weisenstein, P.O. Wennberg, *Geophys. Res. Lett.* 32 (2005) L23807.
- [14] N.J. Livesey, M.D. Fromm, J.W. Waters, G.L. Manney, M.L. Santee, W.G. Read, *J. Geophys. Res.* 109 (2004) D06308.
- [15] I.J. Simpson, S.K. Akagi, B. Barletta, N.J. Blake, Y. Choi, G.S. Diskin, A. Fried, H.E. Fuelberg, S. Meinardi, F.S. Rowland, S.A. Vay, A.J. Weinheimer, P.O. Wennberg, P. Wiebring, A. Wisthaler, M. Yang, R.J. Yokelson, D.R. Blake, *Atmos. Chem. Phys.* 11 (2011) 6445–6463.
- [16] P.M. Solomon, K.B. Jefferts, A.A. Penzias, R.W. Wilson, *Astrophys. J.* 168 (1971) L107–L110.
- [17] B.L. Ulich, E.K. Conklin, *Nature* 248 (1974) 121–122.
- [18] B. Bézar, A. Marten, G. Paubert, *Bull. Am. Astron. Soc.* 25 (1993) 1100.
- [19] H.E. Matthews, T.J. Sears, *Astrophys. J.* 267 (1983) L53–L57.
- [20] S. Cazaux, A.G.G.M. Tielens, C. Ceccarelli, A. Castets, V. Wakelam, E. Caux, B. Parise, D. Teyssier, *Astrophys. J.* 593 (2003) L51–L55.
- [21] M. Agúndez, J. P. Fonfría, J. Cernicharo, J. R. Pardo, M. Guélin, *Astron. Astrophys.* 479 (2008) 493–501.
- [22] R. Mauersberger, C. Henkel, C.M. Walmsley, L.J. Sage, T. Wiklind, *Astron. Astrophys.* 247 (1991) 307–314.
- [23] M. Gerin, F. Combes, G. Wlodarczak, T. Jacq, M. Guélin, P. Encrenaz, C. Laurent, *Astron. Astrophys.* 259 (1992) L35–L38.
- [24] P.F. Goldsmith, R. Krotkov, R.L. Snell, R.D. Brown, P. Godfrey, *Astrophys. J.* 274 (1983) 184–194.
- [25] S.M. Fortman, J. P. McMillan, C.F. Neese, S.K. Randall, A.J. Remijan, T.L. Wilson, F.C. De Lucia, *J. Mol. Spectrosc.* 280 (2012) 11–20.
- [26] H. Ring, H. Edwards, M. Kessler, W. Gordy, *Phys. Rev.* 72 (1947) 1262–1263.
- [27] H.S.P. Müller, B.J. Drouin, J.C. Pearson, *Astron. Astrophys.* 506 (2009) 1487–1499.
- [28] G. Cazzoli, C. Pazzarini, *J. Mol. Spectrosc.* 240 (2006) 153–163.
- [29] S.G. Kukolich, D.J. Ruben, J.H.S. Wang, J.R. Williams, *J. Chem. Phys.* 58 (1973) 3155–3159.
- [30] S.G. Kukolich, *J. Chem. Phys.* 76 (1982) 97–101.
- [31] D. Boucher, J. Burie, J. Demaison, A. Dubrulle, J. Legrand, B. Segard, *J. Mol. Spectrosc.* 64 (1977) 290–294.
- [32] R. Anttila, V.M. Horneman, M. Koivusaari, R. Paso, *J. Mol. Spectrosc.* 157 (1993) 198–207.
- [33] M. Koivusaari, V.-M. Horneman, R. Anttila, *J. Mol. Spectrosc.* 152 (1992) 377–388.
- [34] P. Venkateswarlu, J.G. Baker, W. Gordy, *J. Mol. Spectrosc.* 6 (1961) 215–228.
- [35] A. Bauer, S. Maes, *J. Phys. (Paris)* 30 (1969) 169–180.
- [36] A. Bauer, PhD thesis, Université de Lille, France, 1970.
- [37] R. Bocquet, G. Wlodarczak, A. Bauer, J. Demaison, *J. Mol. Spectrosc.* 127 (1988) 382–389.
- [38] W. Chen, R. Bocquet, D. Boucher, J. Demaison, G. Wlodarczak, *J. Mol. Spectrosc.* 150 (1991) 470–467.
- [39] W. Lafferty, *J. Mol. Spectrosc.* 25 (1968) 359–364.
- [40] J. Cosleou, J.C. López, J.L. Alonso, G. Wlodarczak, J. Demaison, *J. Mol. Spectrosc.* 149 (1991) 242–251.
- [41] M. Šimečková, Š. Urban, U. Fuchs, F. Lewen, G. Winnewisser, I. Morino, K.M.T. Yamada, *J. Mol. Spectrosc.* 226 (2004) 123–136.
- [42] M. Koivusaari, A.-M. Tolonen, R. Paso, J. Schroderus, R. Anttila, *J. Mol. Spectrosc.* 160 (1993) 566–573.
- [43] J.L. Duncan, D.C. McKean, F. Tullini, G.D. Nivellini, J. Perez Peña, *J. Mol. Spectrosc.* 69 (1978) 123–140.
- [44] Y. Mori, T. Nakagawa, K. Kuchitsu, *J. Mol. Spectrosc.* 104 (1984) 388–401.
- [45] A.-M. Tolonen, M. Koivusaari, J. Schroderus, S. Alanko, R. Anttila, *J. Mol. Spectrosc.* 160 (1993) 554–565.
- [46] A. Bauer, S. Maes, *C. R. Acad. Sci., Ser. B* 268 (1969) 1569.
- [47] A. Bauer, *J. Mol. Spectrosc.* 40 (1971) 183–206.
- [48] H.S.P. Müller et al. contribution WG03, presented at the 62nd International Symposium on Molecular Spectroscopy, June 18–22, 2007, Columbus, Ohio, USA; see also <http://molspect.chemistry.ohio-state.edu/symposium.62/symposium/Abstracts/p074.pdf>
- [49] M. Hermann, D.S. Perry, *Phys. Chem. Chem. Phys.* 15 (2013) 9970–9993.
- [50] H.S.P. Müller, P. Pracna, V.-M. Horneman, *J. Mol. Spectrosc.* 216 (2002) 397–407.
- [51] P. Pracna, H.S.P. Müller, S. Klee, V.-M. Horneman, *Mol. Phys.* 102 (2004) 1555–1568.
- [52] P. Pracna, H.S.P. Müller, Š. Urban, V.-M. Horneman, S. Klee, *J. Mol. Spectrosc.* 256 (2009) 152–162.
- [53] P. Pracna, J. Urban, O. Votava, Z. Meltzerová, Š. Urban, V.-M. Horneman, *Mol. Phys.* 109 (2011) 2237–2243.
- [54] K. Yamada, *J. Mol. Spectrosc.* 81 (1980) 139–151.
- [55] F. Hegelund, J. Bendtsen, *J. Mol. Spectrosc.* 124 (1987) 306–316.
- [56] Z. Kisiel, L. Pszczółkowski, B.J. Drouin, C.S. Brauer, S. Yu, J.C. Pearson, *J. Mol. Spectrosc.* 258 (2009) 26–34.
- [57] Z. Kisiel, L. Pszczółkowski, B.J. Drouin, C.S. Brauer, S. Yu, J.C. Pearson, I.R. Medvedev, S. Fortman, C. Neese, *J. Mol. Spectrosc.* 280 (2012) 134–144.
- [58] B.J. Drouin, F.W. Maiwald, J.C. Pearson, *Rev. Sci. Instr.* 76 (2005) 093113
- [59] G. Winnewisser, A.F. Krupnov, M.Y. Tretyakov, M. Liedtke, F. Lewen, A.H. Saleck, R. Schieder, A.P. Shkaev, S.V. Volokhov, *J. Mol. Spectrosc.* 165 (1994) 294–300.
- [60] L.-H. Xu, R.M. Lees, G.T. Crabbe, J.A. Myshrall, H.S.P. Müller, C.P. Endres, O. Baum, F. Lewen, S. Schlemmer, K.M. Menten, B.E. Billinghurst, *J. Chem. Phys.* 137 (2012) 104313.
- [61] L.-H. Xu, R.M. Lees, Y. Hao, H.S.P. Müller, C.P. Endres, F. Lewen, S. Schlemmer, *J. Mol. Spectrosc.* 303 (2014) 1–7.
- [62] A.G. Maki, J.S. Wells, *J. Res. Nat. Inst. Stand. Tech.* 97 (1992) 409–470.
- [63] L.R. Brown, J.S. Margolis, R.H. Norton, B. Stedry, *Appl. Spectrosc.* 37 (1983) 287–292.
- [64] L.S. Rothman, I.E. Gordon, Y. Babikov, A. Barbe, D.C. Benner, et al., *J. Quant. Spectrosc. Radiat. Transfer* 130 (2013) 4–50.
- [65] J. Gadhii, A. Lahrouni, J. Legrand, J. Demaison, *J. Chim. Phys. Phys.-Chim. Biol.* 92 (1995) 1984–1992.
- [66] C.P. Rinsland, S.W. Sharpe, R.L. Sams, *J. Quant. Spectrosc. Radiat. Transfer* 96 (2005) 271–280.
- [67] M. Nishio, P. Paillous, M. Khelif, P. Bruston, F. Raulin, *Spectrochim. Acta* 51A (1995) 617–622.
- [68] C.P. Rinsland, V.M. Devi, D.C. Benner, T.A. Blake, R.L. Sams, L.R. Brown, I. Kleiner, A. Dehayem-Kamadjeu, H.S.P. Müller, R.R. Gamache, D.L. Niles, T. Masiello, *J. Quant. Spectrosc. Radiat. Transfer* 109 (2009) 974–994.
- [69] R. Paso, R. Anttila, M. Koivusaari, *J. Mol. Spectrosc.* 165 (1994) 470–480.
- [70] K. Sarka, L. Nová Střiteská, A. Ceausu-Velcescu, *J. Mol. Spectrosc.* (2014), in press, doi: 10.1016/j.jms.2014.09.003.
- [71] H.M. Pickett, *J. Mol. Spectrosc.* 148 (1991) 371–377.
- [72] B.J. Drouin, H.S.P. Müller, *J. Mol. Spectrosc.* 251 (2008) 1–3.
- [73] J.C. Pearson, H.S.P. Müller, H.M. Pickett, E.A. Cohen, B.J. Drouin, *J. Quant. Spectrosc. Radiat. Transfer* 111 (2010) 1614–1616.
- [74] S. Yu, B.J. Drouin, J.C. Pearson, H.M. Pickett, *Astrophys. J. Suppl. Ser.* 180 (2009) 119–124.
- [75] J. Plíva, L.D. Le, J.W.C. Johns, Z. Lu, R.A. Bernheim, *J. Mol. Spectrosc.*

- 173 (1995) 423–430.
- [76] P. Pracna, J. Urban, O. Votava, Z. Meltzerová, Š. Urban, V.-M. Horneman, B.J. Drouin, *J. Phys. Chem. A* 115 (2011) 1063–1068.
- [77] S. Thorwirth, H.S.P. Müller, G. Winnewisser, *J. Mol. Spectrosc.* 204 (2000) 133–144.
- [78] H. Spahn, H.S.P. Müller, T.F. Giesen, J.-U. Grabow, M.E. Harding, J. Gauss, S. Schlemmer, *Chem. Phys.* 346 (2008) 132–138.
- [79] H.M. Pickett, R.L. Poynter, E.A. Cohen, M.L. Delitsky, J.C. Pearson, H.S.P. Müller, *J. Quant. Spectrosc. Radiat. Transfer* 60 (1998) 883–890.
- [80] N. Jacquinet-Husson, L. Crepeau, R. Armante, C. Boutammine, A. Chédin, et al., *J. Quant. Spectrosc. Radiat. Transfer* 112 (2011) 2395–2445.
- [81] L.R. Brown, J.S. Margolis, *J. Quant. Spectrosc. Radiat. Transfer* 56 (1996) 283–294.
- [82] V. Nemtchinov, K. Sung, P. Varanasi, *J. Quant. Spectrosc. Radiat. Transfer* 83 (2004) 243–265.
- [83] J.M. Colmont, F. Rohart, G. Wlodarczak, J.P. Bouanich, *J. Mol. Spectrosc.* 238 (2006) 98–107.
- [84] M. Baranger, *Phys. Rev.* 112 (1958) 855–865.
- [85] F. Cerceau, F. Raulin, R. Courtin, D. Gautier, *Icarus* 62 (1985) 207–220.
- [86] Y. Koga, S. Kondo, S. Saeki, W.B. Person, *J. Phys. Chem.* 88 (1984) 3152–3157.